A weak Harnack inequality for fractional evolution equations with discontinuous coefficients

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Abstract. We study linear time fractional diffusion equations in divergence form of time order less than one. It is merely assumed that the coefficients are measurable and bounded, and that they satisfy a uniform parabolicity condition. As a main result we establish for nonnegative weak supersolutions of such problems a weak Harnack inequality with optimal critical exponent. The proof relies on new a priori estimates for time fractional problems and uses Moser's iteration technique and an abstract lemma of Bombieri and Giusti, the latter allowing to avoid the rather technically involved approach via *BMO*. As applications of the weak Harnack inequality we establish the strong maximum principle, the continuity of weak solutions at t = 0, and a uniqueness theorem for global bounded weak solutions.

Mathematics Subject Classification (2010): 35R09 (primary); 45K05 (secondary).

1. Introduction and main result

Let T > 0 and Ω be a bounded domain in \mathbb{R}^N . In this paper we are concerned with linear partial integro-differential equations of the form

$$\partial_t^{\alpha}(u-u_0) - \operatorname{div}(A(t,x)Du) = 0, \quad t \in (0,T), \ x \in \Omega.$$
 (1.1)

Here $u_0 = u_0(x)$ is a given initial data for u, $A = (a_{ij})$ is $\mathbb{R}^{N \times N}$ -valued, Du denotes the spatial gradient of u, and ∂_t^{α} stands for the Riemann-Liouville fractional derivation operator with respect to time of order $\alpha \in (0, 1)$; it is defined by

$$\partial_t^{\alpha} v(t, x) = \partial_t \int_0^t g_{1-\alpha}(t-\tau) v(\tau, x) \, d\tau, \quad t > 0, \ x \in \Omega,$$

Work partially supported by the European Community's Human Potential Programme [Evolution Equations for Deterministic and Stochastic Systems], contract code HPRN-CT-2002-00281, and by the Deutsche Forschungsgemeinschaft (DFG), Bonn, Germany.

Received July 23, 2011; accepted December 12, 2011.

where g_{β} denotes the Riemann-Liouville kernel

$$g_{\beta}(t) = \frac{t^{\beta-1}}{\Gamma(\beta)}, \quad t > 0, \ \beta > 0.$$

As to applications, equation (1.1) is a special case of problems arising in mathematical physics when describing dynamic processes in materials with memory, *e.g.* in the theory of heat conduction with memory, see [24] and the references therein. Time fractional diffusion equations are also used to model anomalous diffusion, see *e.g.* [20]. In this context, equations of the type (1.1) are termed *subdiffusion equations* (the time order α lies in (0, 1); in the case $\alpha \in (1, 2)$, which is not considered here, one speaks of *superdiffusion equations*. Time fractional diffusion equations of time order $\alpha \in (0, 1)$ are closely related to a class of Montroll-Weiss continuous time random walk models where the waiting time density behaves as $t^{-\alpha-1}$ for $t \to \infty$, see *e.g.* [14,15,20]. Problems of the type (1.1) are further used to describe diffusion on fractals ([20,25]), and they also appear in mathematical finance, see *e.g.* [27].

Letting $\Omega_T = (0, T) \times \Omega$ we will assume that

(H1) $A \in L_{\infty}(\Omega_T; \mathbb{R}^{N \times N})$, and

$$\sum_{i,j=1}^{N} |a_{ij}(t,x)|^2 \le \Lambda^2, \quad \text{for a.a. } (t,x) \in \Omega_T.$$

(H2) There exists v > 0 such that

 $(A(t, x)\xi|\xi) \ge \nu|\xi|^2$, for a.a. $(t, x) \in \Omega_T$, and all $\xi \in \mathbb{R}^N$.

(H3) $u_0 \in L_2(\Omega)$.

We say that a function u is a *weak solution (subsolution, supersolution)* of (1.1) in Ω_T , if u belongs to the space

$$Z_{\alpha} := \{ v \in L_{\frac{2}{1-\alpha}, w}([0, T]; L_{2}(\Omega)) \cap L_{2}([0, T]; H_{2}^{1}(\Omega)) \text{ such that} \\ g_{1-\alpha} * v \in C([0, T]; L_{2}(\Omega)), \text{ and } (g_{1-\alpha} * v)|_{t=0} = 0 \},$$

and for any nonnegative test function

$$\eta \in \mathring{H}_{2}^{1,1}(\Omega_{T}) := H_{2}^{1}([0,T]; L_{2}(\Omega)) \cap L_{2}([0,T]; \mathring{H}_{2}^{1}(\Omega)) \left(\mathring{H}_{2}^{1}(\Omega) := \overline{C_{0}^{\infty}(\Omega)}^{H_{2}^{1}(\Omega)} \right)$$

with $\eta|_{t=T} = 0$ there holds

$$\int_{0}^{T} \int_{\Omega} \left(-\eta_{t} [g_{1-\alpha} * (u-u_{0})] + (ADu|D\eta) \right) dx dt = (\leq, \geq) 0.$$
 (1.2)

Here $L_{p,w}$ denotes the weak L_p space and $f_1 * f_2$ the convolution on the positive halfline with respect to time, that is $(f_1 * f_2)(t) = \int_0^t f_1(t-\tau) f_2(\tau) d\tau, t \ge 0$.

Weak solutions of (1.1) in the class Z_{α} have been constructed in [37]. Notice also that the function u_0 plays the role of the initial data for u, at least in a weak sense. In case of sufficiently smooth functions u and $g_{1-\alpha} * (u - u_0)$ the condition $(g_{1-\alpha} * u)|_{t=0} = 0$ implies $u|_{t=0} = u_0$, see [37].

To formulate our main result, let B(x, r) denote the open ball with radius r > 0 centered at $x \in \mathbb{R}^N$. By μ_N we mean the Lebesgue measure in \mathbb{R}^N . For $\delta \in (0, 1)$, $t_0 \ge 0, \tau > 0$, and a ball $B(x_0, r)$, define the boxes

$$Q_{-}(t_0, x_0, r) = (t_0, t_0 + \delta \tau r^{2/\alpha}) \times B(x_0, \delta r),$$

$$Q_{+}(t_0, x_0, r) = (t_0 + (2 - \delta)\tau r^{2/\alpha}, t_0 + 2\tau r^{2/\alpha}) \times B(x_0, \delta r).$$

Theorem 1.1. Let $\alpha \in (0, 1)$, T > 0, and $\Omega \subset \mathbb{R}^N$ be a bounded domain. Suppose the assumptions (H1)-(H3) are satisfied. Let further $\delta \in (0, 1)$, $\eta > 1$, and $\tau > 0$ be fixed. Then for any $t_0 \ge 0$ and r > 0 with $t_0 + 2\tau r^{2/\alpha} \le T$, any ball $B(x_0, \eta r) \subset \Omega$, any 0 , and any nonnegative weak supersolution <math>u of (1.1) in $(0, t_0 + 2\tau r^{2/\alpha}) \times B(x_0, \eta r)$ with $u_0 \ge 0$ in $B(x_0, \eta r)$, there holds

$$\left(\frac{1}{\mu_{N+1}(Q_{-}(t_0, x_0, r))} \int_{Q_{-}(t_0, x_0, r)} u^p \, d\mu_{N+1}\right)^{1/p} \le C \underset{Q_{+}(t_0, x_0, r)}{\operatorname{ess \, inf}} u, \quad (1.3)$$

where the constant $C = C(\nu, \Lambda, \delta, \tau, \eta, \alpha, N, p)$.

Theorem 1.1 states that nonnegative weak supersolutions of (1.1) satisfy a weak form of Harnack inequality in the sense that we do not have an estimate for the supremum of u on $Q_{-}(t_0, x_0, r)$ but only an L_p estimate. We also show that the critical exponent $\frac{2+N\alpha}{2+N\alpha-2\alpha}$ is optimal, i.e. the inequality fails to hold for $p \ge \frac{2+N\alpha}{2+N\alpha-2\alpha}$.

Theorem 1.1 can be viewed as the time fractional analogue of the corresponding result in the classical parabolic case $\alpha = 1$, see *e.g.* [19, Theorem 6.18] and [29]. Sending $\alpha \to 1$ in the expression for the critical exponent yields 1 + 2/N, which is the well-known critical exponent for the heat equation. We would like to point out that the statement of Theorem 1.1 remains valid for (appropriately defined) weak supersolutions of (1.1) on $(t_0, t_0 + 2\tau r^{2/\alpha}) \times B(x_0, \eta r)$ which are nonnegative on $(0, t_0 + 2\tau r^{2/\alpha}) \times B(x_0, \eta r)$. To avoid further technicalities we have confined ourselves to consider only weak supersolutions on the full time-interval $(0, t_0 + 2\tau r^{2/\alpha})$. Note that the global positivity assumption cannot be replaced by a local one, as simple examples show, *cf.* [36]. This significant difference to the case $\alpha = 1$ is due to the non-local nature of ∂_t^{α} . The same phenomenon is known for integro-differential operators like $(-\Delta)^{\alpha}$ with $\alpha \in (0, 1)$, see *e.g.* [16].

As a simple consequence of the weak Harnack inequality we derive the strong maximum principle for weak subsolutions of (1.1), see Theorem 5.1 below. The weak maximum principle has been proven in [33], even in a more general setting.

As a further consequence of the weak Harnack inequality we obtain a uniqueness theorem for global bounded weak solutions, see Corollary 5.3 below. It states that any bounded weak solution of (1.1) on $\mathbb{R}_+ \times \mathbb{R}^N$ with $u_0 = 0$ vanishes a.e. on $\mathbb{R}_+ \times \mathbb{R}^N$.

The proof of Theorem 1.1 relies on new a priori estimates for time fractional problems, which are derived by means of the fundamental identity (2.6) (see below) for the regularized fractional derivative. It further uses Moser's iteration technique and an elementary but subtle lemma of Bombieri and Giusti [2], which allows to avoid the rather technically involved approach via BMO-functions. This simplification is already of great significance in the classical parabolic case, see Moser [23] and Saloff-Coste [26].

One of the technical difficulties in deriving the desired estimates in the weak setting is to find an appropriate time regularization of the problem. In the case $\alpha = 1$ this can be achieved by means of Steklov averages in time. In the time fractional case this method does not work anymore, since Steklov average operators and convolution do not commute. It turns out that instead one can use the Yosida approximation of the fractional derivative, which leads to a regularization of the kernel $g_{1-\alpha}$. This method has already been used in [12,30,37], and [33].

We point out that the results obtained in this paper can be generalized to quasilinear equations of the form

$$\partial_t^{\alpha}(u - u_0) - \operatorname{div} a(t, x, u, Du) = b(t, x, u, Du), \ t \in (0, T), \ x \in \Omega,$$
(1.4)

with suitable structure conditions on the functions a and b. This is possible, as also known from the elliptic and the classical parabolic case, since the test function method used in the proof of Theorem 1.1 does not depend so much on the linearity of the differential operator w.r.t. the spatial variables but on a certain nonlinear structure, *cf*. [11,19,29], and [33].

In the literature there exist many papers where equations of the type (1.1), as well as nonlinear or abstract variants of them are studied in a *strong* setting, assuming more smoothness on the coefficients and nonlinearities, see *e.g.* [1,4,7, 10,12,24,34,35]. Concerning the *weak* setting described above one finds only a few results. Existence of weak solutions has been shown in [37] in an abstract setting for a more general class of kernels. Boundedness of weak solutions has been obtained in [33] in the quasilinear case by means of the De Giorgi technique. In the very recent work [31], it is proved that any weak solution of (1.1) is Hölder continuous in the interior of the parabolic domain. The proof is quite involved and relies on the De Giorgi technique and the method of non-local growth lemmata (*cf.* [28]). With the weak Harnack inequality, the present paper establishes another key result towards a De Giorgi-Nash-Moser theory for time fractional evolution equations in divergence form of order $\alpha \in (0, 1)$.

In the classical parabolic case boundedness and the weak (or full) Harnack inequality imply an Hölder estimate for weak solutions, cf. [9, 18, 19, 22]. We also refer to [11] and [21] for the elliptic case. In the present situation one cannot argue anymore as in the classical parabolic case, due to the global positivity assumption in Theorem 1.1. The same problem arises for the fractional Laplacian, see [28]. However, in our case it is possible to establish without much effort at least continuity at t = 0. This is done in Theorem 5.2 in the case $u_0 = 0$. It is shown that in this case any bounded weak solution u of (1.1) is continuous at $(0, x_0)$ for all $x_0 \in \Omega$ and $\lim_{(t,x)\to(0,x_0)} u(t, x) = 0$. Thus for such weak solutions the initial condition $u|_{t=0} = 0$ is satisfied in the classical sense.

We further remark that in the purely time-dependent case, that is for scalar equations of the form

$$\partial_t^{\alpha}(u-u_0) + \sigma u = 0, \quad t \in (0,T),$$

with $\sigma \ge 0$, a weak Harnack inequality with optimal exponent $1/(1 - \alpha)$ has been proven in [32] for nonnegative supersolutions. Recently, the full Harnack inequality for nonnegative solutions has been established in [36]. This, together with the above results, indicates that the full Harnack inequality should also hold for nonnegative solutions of (1.1), which is still an open problem, even in the case $A(t, x) \equiv Id$.

The paper is organized as follows. In Section 2 we collect the basic tools needed for the proof of Theorem 1.1. These include two abstract lemmas on Moser iterations and the lemma of Bombieri and Giusti. We further explain the approximation method for the fractional derivation operator and state the fundamental identity (2.6), which is frequently used in Section 3, where we give the proof of the main result. In Section 4 we show that the critical exponent in Theorem 1.1 is optimal. Finally, Section 5 is devoted to applications of the weak Harnack inequality.

ACKNOWLEDGEMENTS. This paper was initiated while the author was visiting the Technical University Delft (NL) in 2003/2004. The author is greatly indebted to Philippe Clément for many fruitful discussions and valuable suggestions.

2. Preliminaries

2.1. Moser iterations and an abstract lemma of Bombieri and Giusti

Throughout this subsection U_{σ} , $0 < \sigma \leq 1$, will denote a collection of measurable subsets of a fixed finite measure space endowed with a measure μ , such that $U_{\sigma'} \subset U_{\sigma}$ if $\sigma' \leq \sigma$. For $p \in (0, \infty)$ and $0 < \sigma \leq 1$, $L_p(U_{\sigma})$ stands for the Lebesgue space $L_p(U_{\sigma}, d\mu)$ of all μ -measurable functions $f : U_{\sigma} \to \mathbb{R}$ with $|f|_{L_p(U_{\sigma})} := (\int_{U_{\sigma}} |f|^p d\mu)^{1/p} < \infty$.

The following two lemmas are basic to Moser's iteration technique. The arguments in their proofs have been repeatedly used in the literature (see *e.g.* [11,19,21, 22,26,29]), so it is worthwhile to formulate them as lemmas in abstract form, also for future reference. We provide proofs for the sake of completeness.

The first Moser iteration result reads as follows, see also [8, Lemma 2.3].

Lemma 2.1. Let $\kappa > 1$, $\bar{p} \ge 1$, $C \ge 1$, and $\gamma > 0$. Suppose f is a μ -measurable function on U_1 such that

$$|f|_{L_{\beta\kappa}(U_{\sigma'})} \le \left(\frac{C(1+\beta)^{\gamma}}{(\sigma-\sigma')^{\gamma}}\right)^{1/\beta} |f|_{L_{\beta}(U_{\sigma})}, \quad 0 < \sigma' < \sigma \le 1, \ \beta > 0.$$
(2.1)

Then there exist constants $M = M(C, \gamma, \kappa, \bar{p})$ and $\gamma_0 = \gamma_0(\gamma, \kappa)$ such that

$$\operatorname{ess\,sup}_{U_{\delta}} |f| \le \left(\frac{M}{(1-\delta)^{\gamma_0}}\right)^{1/p} |f|_{L_p(U_1)} \quad for \ all \ \delta \in (0, 1), \ p \in (0, \bar{p}].$$

Proof. For q > 0 and $0 < \sigma \le 1$, let

$$\Phi(q,\sigma) = \left(\int_{U_{\sigma}} |f|^q \, d\mu\right)^{1/q}.$$

Let $0 and <math>\delta \in (0, 1)$. Set $p_i = p\kappa^i$, i = 0, 1, ... and define the sequence $\{\sigma_i\}, i = 0, 1, ..., \text{by } \sigma_0 = 1$ and $\sigma_i = 1 - \sum_{j=1}^i 2^{-j}(1-\delta), i = 1, 2, ...;$ observe that $1 = \sigma_0 > \sigma_1 > ... > \sigma_i > \sigma_{i+1} > \delta$ as well as $\sigma_{i-1} - \sigma_i = 2^{-i}(1-\delta), i \ge 1$. Suppose now $n \in \mathbb{N}$. By using (2.1) with $\beta = p_i, i = 0, 1, ..., n-1$, we obtain

$$\begin{split} \Phi(p_{n},\delta) &\leq \Phi(p_{n},\sigma_{n}) = \Phi(p_{n-1}\kappa,\sigma_{n}) \\ &\leq \left(\frac{C(1+p\kappa^{n-1})^{\gamma}}{[2^{-n}(1-\delta)]^{\gamma}}\right)^{\frac{1}{p}\kappa^{-(n-1)}} \Phi(p_{n-1},\sigma_{n-1}) \\ &\leq \left(\frac{C(2\bar{p}\kappa^{n-1})^{\gamma}}{[2^{-n}(1-\delta)]^{\gamma}}\right)^{\frac{1}{p}\kappa^{-(n-1)}} \Phi(p_{n-1},\sigma_{n-1}) \\ &\leq \left(\frac{\tilde{C}(C,\bar{p},\gamma)^{n}\kappa^{\gamma(n-1)}}{(1-\delta)^{\gamma}}\right)^{\frac{1}{p}\kappa^{-(n-1)}} \Phi(p_{n-1},\sigma_{n-1}) \leq \dots \\ &\leq \left(\tilde{C}^{\sum_{j=0}^{n-1}(j+1)\kappa^{-j}}\kappa^{\gamma}\sum_{j=0}^{n-1}j\kappa^{-j}(1-\delta)^{-\gamma}\sum_{j=0}^{n-1}\kappa^{-j}\right)^{1/p} \Phi(p_{0},\sigma_{0}) \\ &\leq \left(\frac{M(C,\bar{p},\gamma,\kappa)}{(1-\delta)^{\frac{\gamma\kappa}{\kappa-1}}}\right)^{1/p} \Phi(p,1). \end{split}$$

We let now *n* tend to ∞ and use the fact that

$$\lim_{n \to \infty} \Phi(p_n, \delta) = \operatorname{ess\,sup}_{U_{\delta}} |f|$$

to get

$$\operatorname{ess\,sup}_{U_{\delta}} |f| \leq \left(\frac{M(C, \bar{p}, \gamma, \kappa)}{(1-\delta)^{\frac{\gamma\kappa}{\kappa-1}}} \right)^{1/p} |f|_{L_{p}(U_{1})}.$$

Hence the proof is complete.

The second Moser iteration result is the following, see also [8, Lemma 2.5].

Lemma 2.2. Assume that $\mu(U_1) \leq 1$. Let $\kappa > 1$, $0 < p_0 < \kappa$, and $C \geq 1$, $\gamma > 0$. Suppose f is a μ -measurable function on U_1 such that

$$|f|_{L_{\beta\kappa}(U_{\sigma'})} \leq \left(\frac{C}{(\sigma - \sigma')^{\gamma}}\right)^{1/\beta} |f|_{L_{\beta}(U_{\sigma})},$$

$$0 < \sigma' < \sigma \leq 1, \ 0 < \beta \leq \frac{p_0}{\kappa} < 1.$$

(2.2)

Then there exist constants $M = M(C, \gamma, \kappa)$ and $\gamma_0 = \gamma_0(\gamma, \kappa)$ such that

$$|f|_{L_{p_0}(U_{\delta})} \leq \left(\frac{M}{(1-\delta)^{\gamma_0}}\right)^{1/p-1/p_0} |f|_{L_p(U_1)} \quad for \ all \ \delta \in (0,1), \ p \in \left(0,\frac{p_0}{\kappa}\right].$$

Proof. Set $p_i = p_0 \kappa^{-i}$, i = 1, 2, ... Given $\delta \in (0, 1)$ we take again the sequence $\{\sigma_i\}, i = 0, 1, 2, ...,$ defined by $\sigma_0 = 1$ and $\sigma_i = 1 - \sum_{j=1}^i 2^{-j} (1 - \delta), i \ge 1$. Suppose now $n \in \mathbb{N}$. By using (2.2) with $\beta = p_i, i = 1, ..., n$, we obtain

$$\Phi(p_{0},\delta) \leq \Phi(p_{0},\sigma_{n}) = \Phi(p_{1}\kappa,\sigma_{n}) \leq \frac{C^{\kappa/p_{0}}}{[2^{-n}(1-\delta)]^{\gamma\kappa/p_{0}}} \Phi(p_{1},\sigma_{n-1})$$

$$\leq \frac{C^{\kappa/p_{0}}}{[2^{-n}(1-\delta)]^{\gamma\kappa/p_{0}}} \frac{C^{\kappa^{2}/p_{0}}}{[2^{-(n-1)}(1-\delta)]^{\gamma\kappa^{2}/p_{0}}} \Phi(p_{2},\sigma_{n-2}) \leq \dots$$

$$\leq \frac{C^{\frac{1}{p_{0}}(\kappa+\kappa^{2}+\ldots+\kappa^{n})}}{2^{-\frac{\gamma}{p_{0}}(n\kappa+(n-1)\kappa^{2}+\ldots+2\kappa^{n-1}+\kappa^{n})}(1-\delta)^{\frac{\gamma}{p_{0}}(\kappa+\kappa^{2}+\ldots+\kappa^{n})}} \Phi(p_{n},\sigma_{0}).$$

Since $p_i = p_0 \kappa^{-i}$, we have

$$\frac{1}{p_0} \sum_{j=1}^n \kappa^j = \frac{\kappa(\kappa^n - 1)}{p_0(\kappa - 1)} = \frac{\kappa}{p_0(\kappa - 1)} \left(\frac{p_0}{p_n} - 1\right) = \frac{\kappa}{\kappa - 1} \left(\frac{1}{p_n} - \frac{1}{p_0}\right).$$

Employing the formula

$$\sum_{j=1}^{n} j\kappa^{j-1} = \frac{1 - (n+1)\kappa^n + n\kappa^{n+1}}{(\kappa - 1)^2}$$

we have further

$$\sum_{j=1}^{n} (n+1-j)\kappa^{j} = (n+1)\sum_{j=1}^{n} \kappa^{j} - \sum_{j=1}^{n} j\kappa^{j}$$
$$= (n+1)\kappa \frac{\kappa^{n}-1}{\kappa-1} - \kappa \frac{1-(n+1)\kappa^{n}+n\kappa^{n+1}}{(\kappa-1)^{2}}$$
$$= \kappa \frac{\kappa^{n+1}-(n+1)\kappa+n}{(\kappa-1)^{2}} \le \frac{\kappa}{(\kappa-1)^{2}}\kappa^{n+1}$$
$$\le \frac{\kappa^{3}}{(\kappa-1)^{3}}(\kappa^{n}-1) \le \frac{\kappa^{3}}{(\kappa-1)^{3}}\left(\frac{p_{0}}{p_{n}}-1\right),$$

which yields

$$\frac{1}{p_0} \sum_{j=1}^n (n+1-j)\kappa^j \le \frac{\kappa^3}{(\kappa-1)^3} \left(\frac{1}{p_n} - \frac{1}{p_0}\right).$$

Therefore

$$\Phi(p_0,\delta) \leq \left[\frac{2^{\frac{\gamma\kappa^3}{(\kappa-1)^3}}C^{\frac{\kappa}{\kappa-1}}}{(1-\delta)^{\frac{\gamma\kappa}{\kappa-1}}}\right]^{\frac{1}{p_n}-\frac{1}{p_0}}\Phi(p_n,\sigma_0).$$

Given $p \in (0, p_0/\kappa]$ there exists $n \ge 2$ such that $p_n . We then have$

$$\frac{1}{p_n} - \frac{1}{p_0} = \frac{\kappa^n - 1}{p_0} \le \frac{\kappa^n + \kappa^{n-1} - \kappa - 1}{p_0} = \frac{(1+\kappa)(\kappa^{n-1} - 1)}{p_0}$$
$$= (1+\kappa)\left(\frac{1}{p_{n-1}} - \frac{1}{p_0}\right) \le (1+\kappa)\left(\frac{1}{p} - \frac{1}{p_0}\right),$$

as well as

$$\Phi(p_n, \sigma_0) = \Phi(p_n, 1) \le \Phi(p, 1),$$

by Hölder's inequality and the assumption $\mu(U_1) \leq 1$. All in all, we obtain

$$\Phi(p_0,\delta) \leq \left[\frac{2^{\frac{\gamma\kappa^3}{(\kappa-1)^3}}C^{\frac{\kappa}{\kappa-1}}}{(1-\delta)^{\frac{\gamma\kappa}{\kappa-1}}}\right]^{(1+\kappa)(\frac{1}{p}-\frac{1}{p_0})}\Phi(p,1),$$

which proves the lemma.

The following abstract lemma is due to Bombieri and Giusti [2]. For a proof we also refer to [26, Lemma 2.2.6] and [8, Lemma 2.6].

Lemma 2.3. Let δ , $\eta \in (0, 1)$, and let γ , *C* be positive constants and $0 < \beta_0 \le \infty$. Suppose *f* is a positive μ -measurable function on U_1 which satisfies the following two conditions:

(i)
$$|f|_{L_{\beta_0}(U_{\sigma'})} \le [C(\sigma - \sigma')^{-\gamma} \mu(U_1)^{-1}]^{1/\beta - 1/\beta_0} |f|_{L_{\beta}(U_{\sigma})}$$

for all σ , σ' , β such that $0 < \delta \le \sigma' < \sigma \le 1$ and $0 < \beta \le \min\{1, \eta\beta_0\}$.

(ii)
$$\mu(\{\log f > \lambda\}) \le C\mu(U_1)\lambda^{-1}$$

for all $\lambda > 0$. Then

$$|f|_{L_{\beta_0}(U_{\delta})} \le M\mu(U_1)^{1/\beta_0},$$

where *M* depends only on δ , η , γ , *C*, and β_0 .

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2.2. The Yosida approximation of the fractional derivation operator

Let $0 < \alpha < 1$, $1 \le p < \infty$, T > 0, and X be a real Banach space. Then the fractional derivation operator B defined by

$$Bu = \frac{d}{dt} (g_{1-\alpha} * u), \quad D(B) = \{ u \in L_p([0, T]; X) : g_{1-\alpha} * u \in {}_0H_p^1([0, T]; X) \},$$

where the zero means vanishing at t = 0, is known to be *m*-accretive in $L_p([0,T];X)$, *cf.* [3, 6], and [12]. Its Yosida approximations B_n , defined by $B_n = nB(n + B)^{-1}$, $n \in \mathbb{N}$, enjoy the property that for any $u \in D(B)$, one has $B_n u \to Bu$ in $L_p([0,T];X)$ as $n \to \infty$. Further, one has the representation

$$B_{n}u = \frac{d}{dt}(g_{1-\alpha,n} * u), \quad u \in L_{p}([0,T];X), \ n \in \mathbb{N},$$
(2.3)

where $g_{1-\alpha,n} = ns_{\alpha,n}$, and $s_{\alpha,n}$ is the unique solution of the scalar-valued Volterra equation

$$s_{\alpha,n}(t) + n(s_{\alpha,n} * g_{\alpha})(t) = 1, \quad t > 0, \ n \in \mathbb{N},$$

see e.g. [30]. Let $h_{\alpha,n} \in L_{1, loc}(\mathbb{R}_+)$ be the resolvent kernel associated with ng_{α} , that is

$$h_{\alpha,n}(t) + n(h_{\alpha,n} * g_{\alpha})(t) = ng_{\alpha}(t), \quad t > 0, \ n \in \mathbb{N}.$$
(2.4)

Convolving (2.4) with $g_{1-\alpha}$ and using $g_{\alpha} * g_{1-\alpha} = 1$, we obtain

$$(g_{1-\alpha} * h_{\alpha,n})(t) + n([g_{1-\alpha} * h_{\alpha,n}] * g_{\alpha})(t) = n, \quad t > 0, \ n \in \mathbb{N}.$$

Hence

$$g_{1-\alpha,n} = n s_{\alpha,n} = g_{1-\alpha} * h_{\alpha,n}, \quad n \in \mathbb{N}.$$
(2.5)

The kernels $g_{1-\alpha,n}$ are nonnegative and nonincreasing for all $n \in \mathbb{N}$, and they belong to $H_1^1([0, T])$, *cf*. [24] and [30]. Note that for any function $f \in L_p([0, T]; X)$, $1 \le p < \infty$, there holds $h_{\alpha,n} * f \to f$ in $L_p([0, T]; X)$ as $n \to \infty$. In fact, setting $u = g_{\alpha} * f$, we have $u \in D(B)$, and

$$B_n u = \frac{d}{dt} (g_{1-\alpha,n} * u) = \frac{d}{dt} (g_{1-\alpha} * g_\alpha * h_{\alpha,n} * f) = h_{\alpha,n} * f \to Bu$$
$$= f \quad \text{in } L_p([0,T];X)$$

as $n \to \infty$. In particular, $g_{1-\alpha,n} \to g_{1-\alpha}$ in $L_1([0, T])$ as $n \to \infty$.

We next state a fundamental identity for integro-differential operators of the form $\frac{d}{dt}(k * u), cf$ also [33]. Suppose $k \in H_1^1([0, T])$ and $H \in C^1(\mathbb{R})$. Then it follows from a straightforward computation that for any sufficiently smooth function u on (0, T) one has for a.a. $t \in (0, T)$,

$$H'(u(t))\frac{d}{dt}(k*u)(t) = \frac{d}{dt}(k*H(u))(t) + (-H(u(t)) + H'(u(t))u(t))k(t) + \int_0^t (H(u(t-s)) - H(u(t)) - H'(u(t))[u(t-s) - u(t)])[-\dot{k}(s)]ds,$$
(2.6)

where k denotes the derivative of k. In particular this identity applies to the Yosida approximations of the fractional derivation operator. We remark that an integrated version of (2.6) can be found in [13, Lemma 18.4.1].

Observe that the last term in (2.6) is nonnegative in case H is convex and k is nonincreasing. This fact is frequently used in the estimates below. However we would like to point out that for the delicate logarithmic estimates in Section 3.3 one really needs the full identity (2.6), even though we work all the time with convex or concave functions. In particular, the crucial fractional differential inequality (3.40) cannot be obtained by using merely convexity inequalities.

The subsequent two lemmas are also obtained by simple algebra.

Lemma 2.4. Let T > 0 and $\alpha \in (0, 1)$. Suppose that $v \in {}_{0}H_{1}^{1}([0, T])$ and $\varphi \in C^{1}([0, T])$. Then

$$\left(g_{\alpha} * (\varphi \dot{v})\right)(t) = \varphi(t)(g_{\alpha} * \dot{v})(t) + \int_{0}^{t} v(\sigma)\partial_{\sigma}\left(g_{\alpha}(t-\sigma)[\varphi(t)-\varphi(\sigma)]\right)d\sigma,$$

a.a. $t \in (0, T).$

If in addition v is nonnegative and φ is nondecreasing there holds

$$\left(g_{\alpha}*(\varphi\dot{v}))(t) \geq \varphi(t)(g_{\alpha}*\dot{v})(t) - \int_{0}^{t} g_{\alpha}(t-\sigma)\dot{\varphi}(\sigma)v(\sigma)\,d\sigma, \ a.a.\ t \in (0,T).\right)$$

Lemma 2.5. Let $T > 0, k \in H_1^1([0, T]), v \in L_1([0, T]), and \varphi \in C^1([0, T]).$ *Then*

$$\varphi(t)\frac{d}{dt}(k*v)(t) = \frac{d}{dt}\left(k*[\varphi v]\right)(t) + \int_0^t \dot{k}(t-\tau)\left(\varphi(t)-\varphi(\tau)\right)v(\tau)\,d\tau, \ a.a.\ t \in (0,T).$$

2.3. An embedding result and a weighted Poincaré inequality

Let T > 0 and Ω be a bounded domain in \mathbb{R}^N . For 1 we define the space

$$V_p := V_p([0, T] \times \Omega) = L_{2p}([0, T]; L_2(\Omega)) \cap L_2([0, T]; H_2^1(\Omega)),$$
(2.7)

endowed with the norm

$$|u|_{V_p([0,T]\times\Omega)} := |u|_{L_{2p}([0,T];L_2(\Omega))} + |Du|_{L_2([0,T];L_2(\Omega))}.$$

Set

$$\kappa := \kappa_p := \frac{2p + N(p-1)}{2 + N(p-1)}$$
(2.8)

with $\kappa_{\infty} = 1 + 2/N$. Then $V_p \hookrightarrow L_{2\kappa}([0, T] \times \Omega)$, and

$$|u|_{L_{2\kappa}([0,T]\times\Omega)} \le C(N,p)|u|_{V_p([0,T]\times\Omega)},$$
(2.9)

for all $u \in V_p \cap L_2([0, T]; \mathring{H}_2^1(\Omega))$. This is a consequence of the Gagliardo-Nirenberg and Hölder's inequality. The case $p = \infty$ is contained, *e.g.*, in [18, page 74 and 75]. The proof given there easily extends to the general case. For a more general embedding result (without proof) we also refer to [33, Section 2].

The following result can be found in [22, Lemma 3], see also [19, Lemma 6.12].

Proposition 2.6. Let $\varphi \in C(\mathbb{R}^N)$ with non-empty compact support of diameter d and assume that $0 \le \varphi \le 1$. Suppose that the domains $\{x \in \mathbb{R}^N : \varphi(x) \ge a\}$ are convex for all $a \le 1$. Then for any function $u \in H_2^1(\mathbb{R}^N)$,

$$\int_{\mathbb{R}^N} \left(u(x) - u_{\varphi} \right)^2 \varphi(x) \, dx \le \frac{2d^2 \mu_N(\operatorname{supp} \varphi)}{|\varphi|_{L_1(\mathbb{R}^N)}} \int_{\mathbb{R}^N} |Du(x)|^2 \varphi(x) \, dx,$$

where

$$u_{\varphi} = \frac{\int_{\mathbb{R}^{N}} u(x)\varphi(x) \, dx}{\int_{\mathbb{R}^{N}} \varphi(x) \, dx}$$

3. Proof of the main result

3.1. The regularized weak formulation, time shifts, and scalings

The following lemma is basic to deriving *a priori* estimates for weak (sub-/super-) solutions of (1.1). It provides an equivalent weak formulation of (1.1) where the singular kernel $g_{1-\alpha}$ is replaced by the more regular kernel $g_{1-\alpha,n}$ ($n \in \mathbb{N}$) given in (2.5). In what follows the kernels $h_n := h_{\alpha,n}, n \in \mathbb{N}$, are defined as in Section 2.2.

Lemma 3.1. Let $\alpha \in (0, 1)$, T > 0, and $\Omega \subset \mathbb{R}^N$ be a bounded domain. Suppose the assumptions (H1)-(H3) are satisfied. Then $u \in Z_{\alpha}$ is a weak solution (subsolution, supersolution) of (1.1) in Ω_T if and only if for any nonnegative function $\psi \in \hat{H}_2^1(\Omega)$ one has

$$\int_{\Omega} \left(\psi \partial_t [g_{1-\alpha,n} * (u-u_0)] + (h_n * [ADu] | D\psi) \right) dx = (\leq, \geq) 0,$$

a.a. $t \in (0, T), n \in \mathbb{N}.$

For a proof we refer to [33, Lemma 3.1], where a more general situation is considered with a slightly different function space for the solution. The proof of Lemma 3.1 is analogous.

Let $u \in Z_{\alpha}$ be a weak supersolution of (1.1) in Ω_T and assume that $u_0 \ge 0$ in Ω . Then Lemma 3.1 and positivity of $g_{1-\alpha,n}$ imply that

$$\int_{\Omega} \left(\psi \partial_t (g_{1-\alpha,n} * u) + (h_n * [ADu] | D\psi) \right) dx \ge 0, \quad \text{a.a. } t \in (0, T), \ n \in \mathbb{N}, \ (3.1)$$

for any nonnegative function $\psi \in \mathring{H}_{2}^{1}(\Omega)$.

Let now $t_1 \in (0, T)$ be fixed. For $t \in (t_1, T)$ we introduce the shifted time $s = t - t_1$ and set $\tilde{f}(s) = f(s + t_1)$, $s \in (0, T - t_1)$, for functions f defined on (t_1, T) . From the decomposition

$$(g_{1-\alpha,n} * u)(t, x) = \int_{t_1}^t g_{1-\alpha,n}(t-\tau)u(\tau, x) d\tau + \int_0^{t_1} g_{1-\alpha,n}(t-\tau)u(\tau, x) d\tau,$$
$$t \in (t_1, T),$$

we then deduce that

$$\partial_t (g_{1-\alpha,n} * u)(t,x) = \partial_s (g_{1-\alpha,n} * \tilde{u})(s,x) + \int_0^{t_1} \dot{g}_{1-\alpha,n}(s+t_1-\tau) u(\tau,x) \, d\tau. \tag{3.2}$$

Assuming in addition that $u \ge 0$ on $(0, t_1) \times \Omega$ it follows from (3.1), (3.2), and the positivity of ψ and of $-\dot{g}_{1-\alpha,n}$ that

$$\int_{\Omega} \left(\psi \,\partial_s (g_{1-\alpha,n} * \tilde{u}) + \left((h_n * [ADu])^{\tilde{-}} | D\psi \right) \right) dx \ge 0,$$

a.a. $s \in (0, T - t_1), n \in \mathbb{N},$ (3.3)

for any nonnegative function $\psi \in \mathring{H}_{2}^{1}(\Omega)$. This relation will be the starting point for all of the estimates below.

We conclude this section with a remark on the scaling properties of equation (1.1). Let $t_0, r > 0$ and $x_0 \in \mathbb{R}^N$. Suppose $u \in Z_\alpha$ is a weak solution (subsolution, supersolution) of (1.1) in $(0, t_0 r^{2/\alpha}) \times B(x_0, r)$. Changing the coordinates according to $s = t/r^{2/\alpha}$ and $y = (x - x_0)/r$ and setting $v(s, y) = u(sr^{2/\alpha}, x_0 + yr), v_0(y) = u_0(x_0 + yr)$, and $\tilde{A}(s, y) = A(sr^{2/\alpha}, x_0 + yr)$, the problem for u is transformed to a problem for v in $(0, t_0) \times B(0, 1)$, namely there holds with $D = D_y$ (also in the weak sense)

$$\partial_s^{\alpha}(v - v_0) - \operatorname{div}\left(\hat{A}(s, y)Dv\right) = (\leq, \geq)0, \quad s \in (0, t_0), \ y \in B(0, 1).$$
(3.4)

3.2. Mean value inequalities

For $\sigma > 0$ we put $\sigma B(x, r) := B(x, \sigma r)$. Recall that μ_N denotes the Lebesgue measure in \mathbb{R}^N .

Theorem 3.2. Let $\alpha \in (0, 1)$, T > 0, and $\Omega \subset \mathbb{R}^N$ be a bounded domain. Suppose the assumptions (H1)-(H3) are satisfied. Let $\eta > 0$ and $\delta \in (0, 1)$ be fixed. Then for any $t_0 \in (0, T]$ and r > 0 with $t_0 - \eta r^{2/\alpha} \ge 0$, any ball $B = B(x_0, r) \subset \Omega$, and any weak supersolution $u \ge \varepsilon > 0$ of (1.1) in $(0, t_0) \times B$ with $u_0 \ge 0$ in B, there holds

$$\operatorname{ess\,sup}_{U_{\sigma'}} u^{-1} \leq \left(\frac{C\mu_{N+1}(U_1)^{-1}}{(\sigma - \sigma')^{\tau_0}}\right)^{1/\gamma} |u^{-1}|_{L_{\gamma}(U_{\sigma})}, \quad \delta \leq \sigma' < \sigma \leq 1, \ \gamma \in (0, 1].$$

Here $U_{\sigma} = (t_0 - \sigma \eta r^{2/\alpha}, t_0) \times \sigma B$, $0 < \sigma \leq 1$, $C = C(\nu, \Lambda, \delta, \eta, \alpha, N)$ and $\tau_0 = \tau_0(\alpha, N)$.

Proof. We may assume that r = 1 and $x_0 = 0$. In fact, in the general case we change coordinates as $t \to t/r^{2/\alpha}$ and $x \to (x - x_0)/r$, thereby transforming the equation to a problem of the same type on $(0, t_0/r^{2/\alpha}) \times B(0, 1)$, cf. Section 3.1.

Fix σ' and σ such that $\delta \leq \sigma' < \sigma \leq 1$ and put $B_1 = \sigma B$. For $\rho \in (0, 1]$ we set $V_{\rho} = U_{\rho\sigma}$. Given $0 < \rho' < \rho \leq 1$, let $t_1 = t_0 - \rho \sigma \eta$ and $t_2 = t_0 - \rho' \sigma \eta$. Then $0 \leq t_1 < t_2 < t_0$. We introduce further the shifted time $s = t - t_1$ and set $\tilde{f}(s) = f(s + t_1), s \in (0, t_0 - t_1)$, for functions f defined on (t_1, t_0) . Since $u_0 \geq 0$ in B and u is a positive weak supersolution of (1.1) in $(0, t_0) \times B$, we have (cf. (3.3))

$$\int_{B} \left(v \partial_{s} (g_{1-\alpha,n} * \tilde{u}) + \left((h_{n} * [ADu])^{\tilde{}} | Dv \right) \right) dx \ge 0,$$

a.a. $s \in (0, t_{0} - t_{1}), n \in \mathbb{N},$ (3.5)

for any nonnegative function $v \in \mathring{H}_{2}^{1}(B)$. For $s \in (0, t_{0} - t_{1})$ we choose the test function $v = \psi^{2} \tilde{u}^{\beta}$ with $\beta < -1$ and $\psi \in C_{0}^{1}(B_{1})$ so that $0 \leq \psi \leq 1, \psi = 1$ in $\rho' B_{1}$, supp $\psi \subset \rho B_{1}$, and $|D\psi| \leq 2/[\sigma(\rho - \rho')]$. By the fundamental identity (2.6) applied to $k = g_{1-\alpha,n}$ and the convex function $H(y) = -(1 + \beta)^{-1}y^{1+\beta}$, y > 0, there holds for a.a. $(s, x) \in (0, t_{0} - t_{1}) \times B$

$$-\tilde{u}^{\beta}\partial_{s}(g_{1-\alpha,n}*\tilde{u}) \geq -\frac{1}{1+\beta}\partial_{s}(g_{1-\alpha,n}*\tilde{u}^{1+\beta}) + \left(\frac{\tilde{u}^{1+\beta}}{1+\beta} - \tilde{u}^{1+\beta}\right)g_{1-\alpha,n}$$

$$= -\frac{1}{1+\beta}\partial_{s}(g_{1-\alpha,n}*\tilde{u}^{1+\beta}) - \frac{\beta}{1+\beta}\tilde{u}^{1+\beta}g_{1-\alpha,n}.$$
(3.6)

We further have

$$Dv = 2\psi D\psi \,\tilde{u}^{\beta} + \beta \psi^2 \tilde{u}^{\beta-1} D\tilde{u}.$$

Using this and (3.6) it follows from (3.5) that for a.a. $s \in (0, t_0 - t_1)$

$$-\frac{1}{1+\beta}\int_{B_1}\psi^2\partial_s(g_{1-\alpha,n}*\tilde{u}^{1+\beta})\,dx+|\beta|\int_{B_1}((h_n*[ADu])\tilde{|}\psi^2\tilde{u}^{\beta-1}D\tilde{u})\,dx$$

$$\leq 2\int_{B_1}\left((h_n*[ADu])\tilde{|}\psi D\psi\,\tilde{u}^\beta\right)\,dx+\frac{\beta}{1+\beta}\int_{B_1}\psi^2\tilde{u}^{1+\beta}g_{1-\alpha,n}\,dx.$$
(3.7)

Next, choose $\varphi \in C^1([0, t_0 - t_1])$ such that $0 \le \varphi \le 1, \varphi = 0$ in $[0, (t_2 - t_1)/2]$, $\varphi = 1$ in $[t_2 - t_1, t_0 - t_1]$, and $0 \le \dot{\varphi} \le 4/(t_2 - t_1)$. Multiplying (3.7) by $-(1+\beta) > 0$ and by $\varphi(s)$, and convolving the resulting inequality with g_α yields

$$\int_{B_{1}} g_{\alpha} * \left(\varphi \partial_{s} (g_{1-\alpha,n} * [\psi^{2} \tilde{u}^{1+\beta}])\right) dx$$

$$+ \beta (1+\beta) g_{\alpha} * \int_{B_{1}} \left((h_{n} * [ADu])^{\tilde{}} |\psi^{2} \tilde{u}^{\beta-1} D \tilde{u} \right) \varphi dx$$

$$\leq 2|1+\beta| g_{\alpha} * \int_{B_{1}} \left((h_{n} * [ADu])^{\tilde{}} |\psi D \psi \tilde{u}^{\beta} \right) \varphi dx$$

$$+ |\beta| g_{\alpha} * \int_{B_{1}} \psi^{2} \tilde{u}^{1+\beta} g_{1-\alpha,n} \varphi dx,$$
(3.8)

for a.a. $s \in (0, t_0 - t_1)$. By Lemma 2.4,

$$\int_{B_1} g_{\alpha} * \left(\varphi \partial_s (g_{1-\alpha,n} * [\psi^2 \tilde{u}^{1+\beta}])\right) dx \ge \int_{B_1} \varphi g_{\alpha} * \left(\partial_s (g_{1-\alpha,n} * [\psi^2 \tilde{u}^{1+\beta}])\right) dx$$

$$- \int_0^s g_{\alpha} (s-\sigma) \dot{\varphi}(\sigma) \left(g_{1-\alpha,n} * \int_{B_1} \psi^2 \tilde{u}^{1+\beta} dx\right)(\sigma) d\sigma.$$
(3.9)

Furthermore, by virtue of

$$g_{1-\alpha,n} * [\psi^2 \tilde{u}^{1+\beta}] \in {}_0H_1^1([0, t_0 - t_1]; L_1(B_1))$$

and $g_{1-\alpha,n} = g_{1-\alpha} * h_n$ as well as $g_{\alpha} * g_{1-\alpha} = 1$ we have

$$g_{\alpha} * \partial_{s}(g_{1-\alpha,n} * [\psi^{2} \tilde{u}^{1+\beta}]) = \partial_{s}(g_{\alpha} * g_{1-\alpha,n} * [\psi^{2} \tilde{u}^{1+\beta}]) = h_{n} * (\psi^{2} \tilde{u}^{1+\beta}).$$
(3.10)

Combining (3.8), (3.9), and (3.10), sending $n \to \infty$, and selecting an appropriate subsequence, if necessary, we thus obtain

$$\begin{split} &\int_{B_1} \varphi \psi^2 \tilde{u}^{1+\beta} \, dx + \beta (1+\beta) \, g_{\alpha} * \int_{B_1} \left(\tilde{A} D \tilde{u} | \psi^2 \tilde{u}^{\beta-1} D \tilde{u} \right) \varphi \, dx \\ &\leq 2 |1+\beta| \, g_{\alpha} * \int_{B_1} \left(\tilde{A} D \tilde{u} | \psi D \psi \, \tilde{u}^{\beta} \right) \varphi \, dx + |\beta| \, g_{\alpha} * \int_{B_1} \psi^2 \tilde{u}^{1+\beta} \, g_{1-\alpha} \varphi \, dx \qquad (3.11) \\ &\quad + \int_0^s g_{\alpha} (s-\sigma) \dot{\varphi}(\sigma) \Big(g_{1-\alpha} * \int_{B_1} \psi^2 \tilde{u}^{1+\beta} \, dx \Big) (\sigma) \, d\sigma, \quad \text{a.a. } s \in (0, t_0 - t_1). \end{split}$$

Put $w = \tilde{u}^{\frac{\beta+1}{2}}$. Then $Dw = \frac{\beta+1}{2}\tilde{u}^{\frac{\beta-1}{2}}D\tilde{u}$. By assumption (H2), we have

$$\beta(1+\beta) g_{\alpha} * \int_{B_{1}} \left(\tilde{A}D\tilde{u}|\psi^{2}\tilde{u}^{\beta-1}D\tilde{u}\right)\varphi dx$$

$$\geq \nu\beta(1+\beta) g_{\alpha} * \int_{B_{1}} \varphi\psi^{2}\tilde{u}^{\beta-1}|D\tilde{u}|^{2} dx \qquad (3.12)$$

$$= \frac{4\nu\beta}{1+\beta} g_{\alpha} * \int_{B_{1}} \varphi\psi^{2}|Dw|^{2} dx.$$

Using (H1) and Young's inequality we may estimate

$$2\left|\left(\tilde{A}D\tilde{u}|\psi D\psi \,\tilde{u}^{\beta}\right)\varphi\right| \leq 2\Lambda\psi|D\psi||D\tilde{u}|\tilde{u}^{\beta}\varphi = 2\Lambda\psi|D\psi||D\tilde{u}|\tilde{u}^{\frac{\beta-1}{2}}\tilde{u}^{\frac{\beta+1}{2}}\varphi$$
$$\leq \frac{\nu|\beta|}{2}\psi^{2}\varphi|D\tilde{u}|^{2}\tilde{u}^{\beta-1} + \frac{2}{\nu|\beta|}\Lambda^{2}|D\psi|^{2}\varphi\tilde{u}^{\beta+1} \quad (3.13)$$
$$= \frac{2\nu|\beta|}{(1+\beta)^{2}}\psi^{2}\varphi|Dw|^{2} + \frac{2}{\nu|\beta|}\Lambda^{2}|D\psi|^{2}\varphi w^{2}.$$

From (3.11), (3.12), and (3.13) we conclude that

$$\int_{B_1} \varphi \psi^2 w^2 dx + \frac{2\nu|\beta|}{|1+\beta|} g_{\alpha} * \int_{B_1} \varphi \psi^2 |Dw|^2 dx \le g_{\alpha} * F, \quad \text{a.a. } s \in (0, t_0 - t_1), \quad (3.14)$$

where

$$F(s) = \frac{2\Lambda^2 |1 + \beta|}{\nu |\beta|} \int_{B_1} |D\psi|^2 \varphi w^2 dx + |\beta|\varphi(s)g_{1-\alpha}(s) \int_{B_1} \psi^2 w^2 dx + \dot{\varphi}(s) \left(g_{1-\alpha} * \int_{B_1} \psi^2 w^2 dx\right)(s) \ge 0, \quad \text{a.a. } s \in (0, t_0 - t_1).$$

We may drop the second term in (3.14), which is nonnegative. By Young's inequality for convolutions and the properties of φ we then infer that for all $p \in (1, 1/(1-\alpha))$

$$\left(\int_{t_2-t_1}^{t_0-t_1} \left(\int_{B_1} [\psi(x)w(s,x)]^2 \, dx\right)^p ds\right)^{1/p} \le |g_{\alpha}|_{L_p([0,t_0-t_1])} \int_0^{t_0-t_1} F(s) \, ds, \quad (3.15)$$

where

$$|g_{\alpha}|_{L_{p}([0,t_{0}-t_{1}])} = \frac{(t_{0}-t_{1})^{\alpha-1+1/p}}{\Gamma(\alpha)[(\alpha-1)p+1]^{1/p}}$$

$$\leq \frac{\eta^{\alpha-1+1/p}}{\Gamma(\alpha)[(\alpha-1)p+1]^{1/p}} =: C_{1}(\alpha, p, \eta).$$
(3.16)

We choose any of these p and fix it.

Returning to (3.14), we may also drop the first term, convolve the resulting inequality with $g_{1-\alpha}$ and evaluate at $s = t_0 - t_1$, thereby obtaining

$$\int_{t_2-t_1}^{t_0-t_1} \int_{B_1} \psi^2 |Dw|^2 \, dx \, ds \le \frac{|1+\beta|}{2\nu|\beta|} \int_0^{t_0-t_1} F(s) \, ds. \tag{3.17}$$

Using

$$\int_{t_2-t_1}^{t_0-t_1} \int_{B_1} |D(\psi w)|^2 \, dx \, ds \le 2 \int_{t_2-t_1}^{t_0-t_1} \int_{B_1} \left(\psi^2 |Dw|^2 + |D\psi|^2 w^2 \right) \, dx \, ds$$

we infer from (3.15)-(3.17) that

$$\begin{aligned} |\psi w|^{2}_{V_{p}([t_{2}-t_{1},t_{0}-t_{1}]\times B_{1})} &\leq 2\left(C_{1}(\alpha, p,\eta) + \frac{|1+\beta|}{\nu|\beta|}\right) \int_{0}^{t_{0}-t_{1}} F(s) \, ds \\ &+ 4\int_{0}^{t_{0}-t_{1}} \int_{B_{1}} |D\psi|^{2} w^{2} \, dx \, ds. \end{aligned}$$
(3.18)

We will next estimate the right-hand side of (3.18). By the assumptions on ψ and φ , and since $|\beta| > 1$, we have

$$\int_0^{t_0-t_1} \int_{B_1} |D\psi|^2 w^2 \, dx \, ds \le \frac{4}{\sigma^2 (\rho - \rho')^2} \int_0^{t_0-t_1} \int_{\rho B_1} w^2 \, dx \, ds$$

and

$$F(s) \leq \left(\frac{8\Lambda^2 |1+\beta|}{\nu\sigma^2(\rho-\rho')^2} + |\beta|g_{1-\alpha}((t_2-t_1)/2)\right) \int_{\rho B_1} w^2 dx + \frac{4}{t_2-t_1} \left(g_{1-\alpha} * \int_{\rho B_1} w^2 dx\right)(s), \quad \text{a.a. } s \in (0, t_0-t_1).$$

Recall that $\sigma \geq \delta > 0$. So we have

$$\begin{split} \int_{0}^{t_{0}-t_{1}} F(s) \, ds &\leq \left(\frac{8\Lambda^{2}|1+\beta|}{\nu\sigma^{2}(\rho-\rho')^{2}} + \frac{2^{\alpha}|\beta|}{\Gamma(1-\alpha)(\rho-\rho')^{\alpha}(\sigma\eta)^{\alpha}} \right) \int_{0}^{t_{0}-t_{1}} \int_{\rho B_{1}} w^{2} \, dx \, ds \\ &+ \frac{4}{(\rho-\rho')\sigma\eta} \int_{0}^{t_{0}-t_{1}} g_{2-\alpha}(t_{0}-t_{1}-\tau) \int_{\rho B_{1}} w(\tau,x)^{2} \, dx \, d\tau \\ &\leq C(\nu,\Lambda,\delta,\eta,\alpha) \frac{1+|1+\beta|}{(\rho-\rho')^{2}} \int_{0}^{t_{0}-t_{1}} \int_{\rho B_{1}} w^{2} \, dx \, ds. \end{split}$$

Combining these estimates and (3.18) yields

$$|\psi w|_{V_p([t_2-t_1,t_0-t_1]\times B_1)} \le C(\nu,\Lambda,\delta,\eta,\alpha,p) \frac{1+|1+\beta|}{\rho-\rho'} |w|_{L_2([0,t_0-t_1]\times \rho B_1)}.$$

We apply next the interpolation inequality (2.9) to the function ψw and make use of $\psi = 1$ in $\rho' B_1$ to deduce that

$$|w|_{L_{2\kappa}([t_2-t_1,t_0-t_1]\times\rho'B_1)} \leq C(\nu,\Lambda,\delta,\eta,\alpha,p,N) \frac{1+|1+\beta|}{\rho-\rho'} |w|_{L_2([0,t_0-t_1]\times\rho B_1)},$$
(3.19)

where the number $\kappa > 1$ is given in (2.8). Since $w = \tilde{u}^{\frac{\beta+1}{2}}$ and by transforming back to the time *t*, we see that (3.19) is equivalent to

$$\left(\int_{V_{\rho'}} u^{-|1+\beta|\kappa} \, d\mu_{N+1}\right)^{\frac{1}{2\kappa}} \leq \frac{\tilde{C}(1+|1+\beta|)}{\rho-\rho'} \left(\int_{V_{\rho}} u^{-|1+\beta|} \, d\mu_{N+1}\right)^{\frac{1}{2}}$$

with $\tilde{C} = \tilde{C}(\nu, \Lambda, \delta, \eta, \alpha, p, N)$. Hence, with $\gamma = |1 + \beta|$,

$$|u^{-1}|_{L_{\gamma\kappa}(V_{\rho'})} \leq \left(\frac{\tilde{C}^2(1+\gamma)^2}{(\rho-\rho')^2}\right)^{1/\gamma} |u^{-1}|_{L_{\gamma}(V_{\rho})}, \quad 0 < \rho' < \rho \leq 1, \ \gamma > 0.$$

Employing the first Moser iteration, Lemma 2.1 (with $\bar{p} = 1$), it follows that there exist constants $M_0 = M_0(\nu, \Lambda, \delta, \eta, \alpha, p, N)$ and $\tau_0 = \tau_0(\kappa)$ such that

$$\operatorname{ess\,sup}_{V_{\theta}} u^{-1} \leq \left(\frac{M_0}{(1-\theta)^{\tau_0}}\right)^{1/\gamma} |u^{-1}|_{L_{\gamma}(V_1)} \quad \text{for all } \theta \in (0,1), \ \gamma \in (0,1].$$

Thus if we take $\theta = \sigma' / \sigma$ and notice that

$$\frac{1}{1-\theta} = \frac{\sigma}{\sigma - \sigma'} \le \frac{1}{\sigma - \sigma'},$$

we obtain

$$\operatorname{ess\,sup}_{U_{\sigma'}} u^{-1} \leq \left(\frac{M_0}{(\sigma - \sigma')^{\tau_0}}\right)^{1/\gamma} |u^{-1}|_{L_{\gamma}(U_{\sigma})}, \quad \gamma \in (0, 1].$$

Hence the proof is complete.

We put

$$\tilde{\kappa} := \kappa_{1/(1-\alpha)} = \frac{2+N\alpha}{2+N\alpha-2\alpha}.$$

Theorem 3.3. Let $\alpha \in (0, 1)$, T > 0, and $\Omega \subset \mathbb{R}^N$ be a bounded domain. Suppose the assumptions (H1)-(H3) are satisfied. Let $\eta > 0$ and $\delta \in (0, 1)$ be fixed. Then for any $t_0 \in [0, T)$ and r > 0 with $t_0 + \eta r^{2/\alpha} \leq T$, any ball $B = B(x_0, r) \subset \Omega$, any $p_0 \in (0, \tilde{\kappa})$, and any nonnegative weak supersolution u of (1.1) in $(0, t_0 + \eta r^{2/\alpha}) \times B$ with $u_0 \geq 0$ in B, there holds

$$\begin{split} |u|_{L_{p_0}(U'_{\sigma'})} &\leq \left(\frac{C\mu_{N+1}(U'_1)^{-1}}{(\sigma-\sigma')^{\tau_0}}\right)^{1/\gamma-1/p_0} |u|_{L_{\gamma}(U'_{\sigma})},\\ \delta &\leq \sigma' < \sigma \leq 1, \ 0 < \gamma \leq p_0/\tilde{\kappa}. \end{split}$$

Here $U'_{\sigma} = (t_0, t_0 + \sigma \eta r^{2/\alpha}) \times \sigma B$, $C = C(\nu, \Lambda, \delta, \eta, \alpha, N, p_0)$, and $\tau_0 = \tau_0(\alpha, N)$.

Proof. We proceed similarly as in the previous proof. Without restriction of generality we may assume that $p_0 > 1$ and r = 1. By replacing u with $u + \varepsilon$ and u_0 with $u_0 + \varepsilon$ and eventually letting $\varepsilon \to 0+$ we may further assume that u is bounded away from zero.

Fix σ' , σ such that $\delta \leq \sigma' < \sigma \leq 1$ and put $B_1 = \sigma B$. For $\rho \in (0, 1]$ we set $V'_{\rho} = U'_{\rho\sigma}$. Given $0 < \rho' < \rho \leq 1$, let $t_1 = t_0 + \rho' \sigma \eta$ and $t_2 = t_0 + \rho \sigma \eta$, so $0 \leq t_0 < t_1 < t_2$. We shift the time by means of $s = t - t_0$ and set $\tilde{f}(s) = f(s + t_0)$, $s \in (0, t_2 - t_0)$, for functions f defined on (t_0, t_2) .

We then repeat the first steps of the preceding proof, the only difference being that now we take $\beta \in (-1, 0)$. Note that, as a consequence of this, (3.6) simplifies to

$$-\tilde{u}^{\beta}\partial_{s}(g_{1-\alpha,n}*\tilde{u}) \geq -\frac{1}{1+\beta}\partial_{s}(g_{1-\alpha,n}*\tilde{u}^{1+\beta}), \quad \text{a.a.} (s,x) \in (0, t_{2}-t_{0}) \times B,$$

hence we obtain with $\psi \in C_0^1(B_1)$ as above

$$-\frac{1}{1+\beta}\int_{B_1}\psi^2\partial_s(g_{1-\alpha,n}*\tilde{u}^{1+\beta})\,dx$$

+ $|\beta|\int_{B_1}\left((h_n*[ADu])\tilde{|}\psi^2\tilde{u}^{\beta-1}D\tilde{u}\right)dx$ (3.20)
 $\leq 2\int_{B_1}\left((h_n*[ADu])\tilde{|}\psi D\psi\,\tilde{u}^\beta\right)dx, \quad \text{a.a. } s \in (0, t_2-t_0).$

Next, choose $\varphi \in C^1([0, t_2 - t_0])$ such that $0 \le \varphi \le 1, \varphi = 1$ in $[0, t_1 - t_0], \varphi = 0$ in $[t_1 - t_0 + (t_2 - t_1)/2, t_2 - t_0]$, and $0 \le -\dot{\varphi} \le 4/(t_2 - t_1)$. Multiplying (3.20) by $1 + \beta > 0$ and by $\varphi(s)$, and applying Lemma 2.5 to the first term gives

$$-\int_{B_{1}}\partial_{s}(g_{1-\alpha,n}*[\varphi\psi^{2}\tilde{u}^{1+\beta}])dx+|\beta|(1+\beta)\int_{B_{1}}(\tilde{A}D\tilde{u}|\psi^{2}\tilde{u}^{\beta-1}D\tilde{u})\varphi dx$$

$$\leq \int_{0}^{s}\dot{g}_{1-\alpha,n}(s-\sigma)(\varphi(s)-\varphi(\sigma))\left(\int_{B_{1}}\psi^{2}\tilde{u}^{1+\beta}dx\right)(\sigma)d\sigma \qquad (3.21)$$

$$+2(1+\beta)\int_{B_{1}}(\tilde{A}D\tilde{u}|\psi D\psi\,\tilde{u}^{\beta})\varphi dx+\mathcal{R}_{n}(s), \quad \text{a.a. } s \in (0, t_{2}-t_{0}),$$

where

$$\mathcal{R}_n(s) = -|\beta|(1+\beta) \int_{B_1} \left((h_n * [ADu])^{\tilde{}} - \tilde{A}D\tilde{u} | \psi^2 \tilde{u}^{\beta-1} D\tilde{u} \right) \varphi \, dx$$
$$+ 2(1+\beta) \int_{B_1} \left((h_n * [ADu])^{\tilde{}} - \tilde{A}D\tilde{u} | \psi D\psi \, \tilde{u}^\beta \right) \varphi \, dx, \quad \text{a.a. } s \in (0, t_2 - t_0).$$

We set again $w = \tilde{u}^{\frac{\beta+1}{2}}$ and estimate exactly as in the preceding proof, using (H1), (H3) and (3.13), to the result

$$-\int_{B_{1}}\partial_{s}(g_{1-\alpha,n}*[\varphi\psi^{2}w^{2}])dx + \frac{2\nu|\beta|}{1+\beta}\int_{B_{1}}\varphi\psi^{2}|Dw|^{2}dx$$

$$\leq \int_{0}^{s}\dot{g}_{1-\alpha,n}(s-\sigma)(\varphi(s)-\varphi(\sigma))(\int_{B_{1}}\psi^{2}w^{2}dx)(\sigma)d\sigma \qquad (3.22)$$

$$+\frac{2\Lambda^{2}(1+\beta)}{\nu|\beta|}\int_{B_{1}}|D\psi|^{2}\varphi w^{2}dx + \mathcal{R}_{n}(s), \quad \text{a.a. } s \in (0, t_{2}-t_{0}).$$

Recall that $g_{1-\alpha,n} = g_{1-\alpha} * h_n$. Putting

$$W(s) = \int_{B_1} \varphi(s)\psi(x)^2 w(s, x)^2 dx$$

and denoting the right-hand side of (3.22) by $F_n(s)$, it follows from (3.22) that

$$G_n(s) := \partial_s^{\alpha}(h_n * W)(s) + F_n(s) \ge 0,$$
 a.a. $s \in (0, t_2 - t_0)$

By (3.10) and positivity of h_n , we have

$$0 \le h_n * W = g_\alpha * \partial_s^\alpha (h_n * W) \le g_\alpha * G_n + g_\alpha * [-F_n(s)]_+$$

a.e. in $(0, t_2 - t_0)$, where $[y]_+$ stands for the positive part of $y \in \mathbb{R}$. For any $p \in (1, 1/(1 - \alpha))$ and any $t_* \in [t_2 - t_0 - (t_2 - t_1)/4, t_2 - t_0]$ we thus obtain by Young's inequality

$$|h_n * W|_{L_p([0,t_*])} \le |g_\alpha|_{L_p([0,t_*])} (|G_n|_{L_1([0,t_*])} + |[-F_n]_+|_{L_1([0,t_*])}).$$
(3.23)

Since $t_* \leq t_2 - t_0 \leq \eta$, we have $|g_{\alpha}|_{L_p([0,t_*])} \leq C_1(\alpha, p, \eta)$ with the same constant as in (3.16). By positivity of G_n ,

$$|G_n|_{L_1([0,t_*])} = (g_{1-\alpha,n} * W)(t_*) + \int_0^{t_*} F_n(s) \, ds.$$

Observe that $\mathcal{R}_n \to 0$ in $L_1([0, t_2 - t_0])$ as $n \to \infty$. Hence $|[-F_n]_+|_{L_1([0, t_*])} \to 0$ as $n \to \infty$. Further,

$$\begin{split} &\int_0^{t_*} \int_0^s \dot{g}_{1-\alpha,n}(s-\sigma) \big(\varphi(s) - \varphi(\sigma)\big) \left(\int_{B_1} \psi^2 w^2 \, dx\right)(\sigma) \, d\sigma \, ds \\ &= \int_0^{t_*} g_{1-\alpha,n}(t_* - \sigma) \big(\varphi(t_*) - \varphi(\sigma)\big) \left(\int_{B_1} \psi^2 w^2 \, dx\right)(\sigma) \, d\sigma \\ &- \int_0^{t_*} \dot{\varphi}(s) \int_0^s g_{1-\alpha,n}(s-\sigma) \left(\int_{B_1} \psi^2 w^2 \, dx\right)(\sigma) \, d\sigma \, ds \\ &\leq - \int_0^{t_*} \dot{\varphi}(s) \int_0^s g_{1-\alpha,n}(s-\sigma) \left(\int_{B_1} \psi^2 w^2 \, dx\right)(\sigma) \, d\sigma \, ds, \end{split}$$

since φ is nonincreasing. We also know that $g_{1-\alpha,n} * W \to g_{1-\alpha} * W$ in $L_1([0, t_2 - t_0])$. Hence we can fix some $t_* \in [t_2 - t_0 - (t_2 - t_1)/4, t_2 - t_0]$ such that for some subsequence $(g_{1-\alpha,n_k} * W)(t_*) \to (g_{1-\alpha} * W)(t_*)$ as $k \to \infty$. Sending $k \to \infty$ it follows then from (3.23), the preceding estimates, and from $\varphi = 1$ in $[0, t_1 - t_0]$ that

$$\left(\int_{0}^{t_{1}-t_{0}} \left(\int_{B_{1}} [\psi(x)w(s,x)]^{2} dx\right)^{p} ds\right)^{1/p} \le C_{1}(\alpha, p, \eta) \left((g_{1-\alpha} * W)(t_{*}) + |F|_{L_{1}([0,t_{2}-t_{0}])}\right),$$
(3.24)

with

$$F(s) = \frac{2\Lambda^2(1+\beta)}{\nu|\beta|} \int_{B_1} |D\psi|^2 \varphi w^2 \, dx - \dot{\varphi}(s) \big(g_{1-\alpha} * \int_{B_1} \psi^2 w^2 \, dx\big)(s).$$

On the other hand, we can integrate (3.22) over $(0, t_*)$ and take the limit as $k \to \infty$ for the same subsequence as before, thereby getting

$$\int_{0}^{t_{1}-t_{0}} \int_{B_{1}} \psi^{2} |Dw|^{2} dx \, ds \leq \frac{1+\beta}{2\nu|\beta|} \left((g_{1-\alpha} * W)(t_{*}) + |F|_{L_{1}([0,t_{2}-t_{0}])} \right). \quad (3.25)$$

Arguing as above (*cf.* the lines before (3.18)), we conclude from (3.24) and (3.25) that

$$\begin{aligned} |\psi w|_{V_{p}([0,t_{1}-t_{0}]\times B_{1})}^{2} &\leq 4 \int_{0}^{t_{2}-t_{0}} \int_{B_{1}} |D\psi|^{2} w^{2} \, dx \, ds \\ &+ 2 \left(C_{1}(\alpha, \, p, \, \eta) + \frac{1+\beta}{\nu|\beta|} \right) \left((g_{1-\alpha} * W)(t_{*}) + |F|_{L_{1}([0,t_{2}-t_{0}])} \right). \end{aligned}$$

$$(3.26)$$

Since $\varphi = 0$ in $[t_1 - t_0 + (t_2 - t_1)/2, t_2 - t_0]$ and $t_* \in [t_2 - t_0 - (t_2 - t_1)/4, t_2 - t_0]$, we have

$$(g_{1-\alpha} * W)(t_*) \leq g_{1-\alpha} ((t_2 - t_1)/4) \int_0^{t_2 - t_0} \int_{B_1} \psi^2 w^2 \, dx \, ds$$

= $\frac{4^{\alpha}}{\Gamma(1-\alpha)(\rho - \rho')^{\alpha}(\sigma \eta)^{\alpha}} \int_0^{t_2 - t_0} \int_{\rho B_1} w^2 \, dx \, ds.$

Further,

$$\int_0^{t_2-t_0} \int_{B_1} |D\psi|^2 w^2 \, dx \, ds \le \frac{4}{\sigma^2 (\rho - \rho')^2} \int_0^{t_2-t_0} \int_{\rho B_1} w^2 \, dx \, ds.$$

The term $|F|_{L_1([0,t_2-t_0])}$ is estimated similarly as in the proof of Theorem 3.2 (*cf.* the lines that follow (3.18)). We obtain

$$|F|_{L_1([0,t_2-t_0])} \le C(\nu,\Lambda,\delta,\eta,\alpha) \frac{1+(1+\beta)}{|\beta|(\rho-\rho')^2} \int_0^{t_2-t_0} \int_{\rho B_1} w^2 \, dx \, ds.$$

Notice the additional factor $|\beta|$ in the denominator. Combining these estimates we deduce from (3.26) that

$$|\psi w|_{V_p([0,t_1-t_0]\times B_1)} \le C(\nu,\Lambda,\delta,\eta,\alpha,p) \frac{1+(1+\beta)}{|\beta|(\rho-\rho')} |w|_{L_2([0,t_2-t_0]\times\rho B_1)}.$$

By the interpolation inequality (2.9) and since $\psi = 1$ in $\rho' B_1$, this implies for all $\beta \in (-1, 0)$

$$|w|_{L_{2\kappa}([0,t_1-t_0]\times\rho'B_1)} \le C(\nu,\Lambda,\delta,\eta,\alpha,p,N) \frac{1+|1+\beta|}{|\beta|(\rho-\rho')} |w|_{L_2([0,t_2-t_0]\times\rho B_1)},$$
(3.27)

where

$$\kappa = \kappa_p = \frac{2p + N(p-1)}{2 + N(p-1)} \in (1, \tilde{\kappa}).$$

We now fix $1 such that <math>\kappa_p = (p_0 + \tilde{\kappa})/2$. This is possible because $\kappa_p \nearrow \tilde{\kappa}$ as $p \nearrow 1/(1-\alpha)$.

Next, we set $\gamma = 1 + \beta \in (0, 1)$ and transform back to *u* to get

$$|u|_{L_{\gamma\kappa}(V'_{\rho'},d\mu)} \le \left(\frac{\tilde{C}}{(\rho-\rho')^2}\right)^{1/\gamma} |u|_{L_{\gamma}(V'_{\rho},d\mu)},$$

$$0 < \rho' < \rho \le 1, \ 0 < \gamma \le p_0/\kappa.$$
(3.28)

Here, $\mu = (\eta \omega_N)^{-1} \mu_{N+1}$, ω_N the volume of the unit ball in \mathbb{R}^N , and $\tilde{C} = \tilde{C}(\nu, \Lambda, \delta, \eta, \alpha, N, p_0)$ is independent of $\gamma \in (0, p_0/\kappa]$, since $|\beta|$ is bounded away from zero. Note that $\mu(V'_1) \leq 1$.

Finally, we employ the second Moser iteration scheme, Lemma 2.2, to conclude from (3.28) that there are constants $M_0 = M_0(\nu, \Lambda, \delta, \eta, \alpha, N, p_0)$ and $\tau_0 = \tau_0(\kappa)$ such that

$$|u|_{L_{p_0}(V'_{\theta},d\mu)} \leq \left(\frac{M_0}{(1-\theta)^{\tau_0}}\right)^{1/\gamma - 1/p_0} |u|_{L_{\gamma}(V'_1,d\mu)},$$

$$0 < \theta < 1, \ 0 < \gamma \leq p_0/\kappa.$$
(3.29)

If we take $\theta = \sigma'/\sigma$ and translate (3.29) back to the measure μ_{N+1} , we obtain

$$|u|_{L_{p_0}(U'_{\sigma'})} \le \left(\frac{M_0(\eta\omega_N)^{-1}}{(\sigma - \sigma')^{\tau_0}}\right)^{1/\gamma - 1/p_0} |u|_{L_{\gamma}(U'_{\sigma})}, \quad 0 < \gamma \le p_0/\kappa.$$
(3.30)

Since $\kappa < \tilde{\kappa}$, (3.30) holds in particular for all $\gamma \in (0, p_0/\tilde{\kappa}]$. This finishes the proof.

3.3. Logarithmic estimates

Theorem 3.4. Let $\alpha \in (0, 1)$, T > 0, and $\Omega \subset \mathbb{R}^N$ be a bounded domain. Suppose the assumptions (H1)–(H3) are satisfied. Let $\tau > 0$ and δ , $\eta \in (0, 1)$ be fixed. Then for any $t_0 \ge 0$ and r > 0 with $t_0 + \tau r^{2/\alpha} \le T$, any ball $B = B(x_0, r) \subset \Omega$, and

any weak supersolution $u \ge \varepsilon > 0$ of (1.1) in $(0, t_0 + \tau r^{2/\alpha}) \times B$ with $u_0 \ge 0$ in *B*, there is a constant c = c(u) such that

 $\mu_{N+1}\big(\{(t,x)\in K_-: \log u(t,x) > c+\lambda\}\big) \le Cr^{2/\alpha}\mu_N(B)\lambda^{-1}, \quad \lambda > 0, \ (3.31)$

and

$$\mu_{N+1}\big(\{(t,x)\in K_+: \log u(t,x) < c-\lambda\}\big) \le Cr^{2/\alpha}\mu_N(B)\lambda^{-1}, \quad \lambda > 0, \quad (3.32)$$

- - -

where $K_{-} = (t_0, t_0 + \eta \tau r^{2/\alpha}) \times \delta B$ and $K_{+} = (t_0 + \eta \tau r^{2/\alpha}, t_0 + \tau r^{2/\alpha}) \times \delta B$. Here the constant C depends only on δ , η , τ , N, α , ν , and Λ .

Proof. Since $u_0 \ge 0$ in *B* and *u* is a positive weak supersolution we may assume without loss of generality that $u_0 = 0$ and $t_0 = 0$. In fact, in the case $t_0 > 0$ we shift the time as $t \to t - t_0$, thereby obtaining an inequality of the same type on the time-interval $J := [0, \tau r^{2/\alpha}]$. Observe that the property $g_{1-\alpha} * u \in C([0, t_0 + \tau r^{2/\alpha}]; L_2(B))$ implies $g_{1-\alpha} * \tilde{u} \in C(J; L_2(B))$ for the shifted function $\tilde{u}(s, x) = u(s + t_0, x)$. So we have

$$\int_{B} \left(v \partial_t (g_{1-\alpha,n} * u) + (h_n * [ADu] | Dv) \right) dx \ge 0, \quad \text{a.a. } t \in J, \ n \in \mathbb{N}, \quad (3.33)$$

for any nonnegative test function $v \in \mathring{H}_{2}^{1}(B)$.

For $t \in J$ we choose the test function $v = \psi^2 u^{-1}$ with $\psi \in C_0^1(B)$ such that supp $\psi \subset B$, $\psi = 1$ in δB , $0 \le \psi \le 1$, $|D\psi| \le 2/[(1 - \delta)r]$ and the domains $\{x \in B : \psi(x)^2 \ge b\}$ are convex for all $b \le 1$. We have

$$Dv = 2\psi D\psi u^{-1} - \psi^2 u^{-2} Du,$$

so that by substitution into (3.33) we obtain for a.a. $t \in J$

$$-\int_{B}\psi^{2}u^{-1}\partial_{t}(g_{1-\alpha,n}*u)\,dx + \int_{B}\left(ADu|u^{-2}Du\right)\psi^{2}\,dx$$

$$\leq 2\int_{B}\left(ADu|u^{-1}\psi D\psi\right)dx + \mathcal{R}_{n}(t),$$
(3.34)

where

$$\mathcal{R}_n(t) = \int_B \left(h_n * [ADu] - ADu | Dv \right) dx$$

By (H1) and Young's inequality,

$$|2(ADu|u^{-1}\psi D\psi)| \le 2\Lambda\psi |D\psi| |Du|u^{-1} \le \frac{\nu}{2}\psi^2 |Du|^2 u^{-2} + \frac{2}{\nu}\Lambda^2 |D\psi|^2.$$

Using this, (H2) and $|D\psi| \le 2/[(1-\delta)r]$, we infer from (3.34) that for a.a. $t \in J$

$$-\int_{B} \psi^{2} u^{-1} \partial_{t} (g_{1-\alpha,n} * u) dx + \frac{\nu}{2} \int_{B} |Du|^{2} u^{-2} \psi^{2} dx \leq \frac{8\Lambda^{2} \mu_{N}(B)}{\nu(1-\delta)^{2} r^{2}} + \mathcal{R}_{n}(t).$$
(3.35)

Setting $w = \log u$ we have $Dw = u^{-1}Du$. The weighted Poincaré inequality of Proposition 2.6 with weight ψ^2 yields

$$\int_{B} (w - W)^{2} \psi^{2} dx \leq \frac{8r^{2} \mu_{N}(B)}{\int_{B} \psi^{2} dx} \int_{B} |Dw|^{2} \psi^{2} dx, \quad \text{a.a. } t \in J,$$
(3.36)

where

$$W(t) = \frac{\int_B w(t, x)\psi(x)^2 dx}{\int_B \psi(x)^2 dx}, \quad \text{a.a. } t \in J.$$

From (3.35) and (3.36) we deduce that

$$-\int_{B} \psi^{2} u^{-1} \partial_{t} (g_{1-\alpha,n} * u) \, dx + \frac{\nu \int_{B} \psi^{2} dx}{16r^{2} \mu_{N}(B)} \int_{B} (w - W)^{2} \psi^{2} dx \leq \frac{8\Lambda^{2} \mu_{N}(B)}{\nu(1-\delta)^{2}r^{2}} + \mathcal{R}_{n}(t),$$

which in turn implies

$$\frac{-\int_{B} \psi^{2} u^{-1} \partial_{t} (g_{1-\alpha,n} * u) dx}{\int_{B} \psi^{2} dx} + \frac{\nu}{16r^{2} \mu_{N}(B)} \int_{\delta B} (w - W)^{2} dx \leq \frac{C_{1}}{r^{2}} + S_{n}(t),$$
(3.37)

for a.a. $t \in J$, with some constant $C_1 = C_1(\delta, N, \nu, \Lambda)$ and $S_n(t) = \mathcal{R}_n(t) / \int_B \psi^2 dx$.

The fundamental identity (2.6) with $H(y) = -\log y$ reads (with the spatial variable x being suppressed)

$$-u^{-1}\partial_t(g_{1-\alpha,n}*u) = -\partial_t(g_{1-\alpha,n}*\log u) + (\log u - 1)g_{1-\alpha,n}(t) + \int_0^t \left(-\log u(t-s) + \log u(t) + \frac{u(t-s) - u(t)}{u(t)}\right) [-\dot{g}_{1-\alpha,n}(s)] ds.$$

In terms of $w = \log u$ this means that

$$-u^{-1}\partial_t(g_{1-\alpha,n} * u) = -\partial_t(g_{1-\alpha,n} * w) + (w-1)g_{1-\alpha,n}(t) + \int_0^t \Psi(w(t-s) - w(t))[-\dot{g}_{1-\alpha,n}(s)] ds,$$
(3.38)

where $\Psi(y) = e^y - 1 - y$. Since Ψ is convex, it follows from Jensen's inequality that

$$\frac{\int_B \psi^2 \Psi \big(w(t-s,x) - w(t,x) \big) \, dx}{\int_B \psi^2 dx} \ge \Psi \left(\frac{\int_B \psi^2 \big(w(t-s,x) - w(t,x) \big) \, dx}{\int_B \psi^2 dx} \right).$$

Using this and (3.38) we obtain

$$\frac{-\int_{B} \psi^{2} u^{-1} \partial_{t}(g_{1-\alpha,n} * u) \, dx}{\int_{B} \psi^{2} dx} \ge -\partial_{t}(g_{1-\alpha,n} * W) + (W-1)g_{1-\alpha,n}(t) + \int_{0}^{t} \Psi (W(t-s) - W(t))[-\dot{g}_{1-\alpha,n}(s)] \, ds \qquad (3.39)$$
$$= -e^{-W} \partial_{t}(g_{1-\alpha,n} * e^{W}),$$

where the last equals sign holds again by (3.38) with *u* replaced by e^W . From (3.37) and (3.39) we conclude that

$$\frac{\nu}{16r^2\mu_N(B)}\int_{\delta B} (w-W)^2 dx \le e^{-W}\partial_t (g_{1-\alpha,n} * e^W) + \frac{C_1}{r^2} + S_n(t), \text{ a.a. } t \in J.$$
(3.40)

We choose

$$c(u) = \log\left(\frac{(g_{1-\alpha} * e^W)(\eta \tau r^{2/\alpha})}{g_{2-\alpha}(\eta \tau r^{2/\alpha})}\right).$$
(3.41)

This definition makes sense, since $g_{1-\alpha} * e^W \in C(J)$. The latter is a consequence of $g_{1-\alpha} * u \in C(J; L_2(B))$ and

$$e^{W(t)} \le \frac{\int_B u(t, x)\psi(x)^2 dx}{\int_B \psi(x)^2 dx}, \quad \text{a.a. } t \in J,$$

where we apply again Jensen's inequality.

To prove (3.31) and (3.32), one of the key ideas is to use the inequalities

$$\mu_{N+1}(\{(t, x) \in K_{-} : w(t, x) > c(u) + \lambda\})$$

$$\leq \mu_{N+1}(\{(t, x) \in K_{-} : w(t, x) > c(u) + \lambda \text{ and } W(t) \leq c(u) + \lambda/2\}) \quad (3.42)$$

$$+ \mu_{N+1}(\{(t, x) \in K_{-} : W(t) > c(u) + \lambda/2\}) =: I_{1} + I_{2}, \quad \lambda > 0,$$

$$\mu_{N+1}(\{(t, x) \in K_{+} : w(t, x) < c(u) - \lambda\})$$

$$\leq \mu_{N+1}(\{(t, x) \in K_{+} : w(t, x) < c(u) - \lambda \text{ and } W(t) \geq c(u) - \lambda/2\}) \quad (3.43)$$

$$+ \mu_{N+1}(\{(t, x) \in K_{+} : W(t) < c(u) - \lambda/2\}) =: I_{3} + I_{4}, \quad \lambda > 0,$$

and to estimate each of the four terms I_i separately.

We begin with the estimates for W. To estimate I_2 and I_4 we adopt some of the ideas developed in [32]. We set $J_- := (0, \eta \tau r^{2/\alpha}), J_+ := (\eta \tau r^{2/\alpha}, \tau r^{2/\alpha})$, and introduce for $\lambda > 0$ the sets $J_-(\lambda) := \{t \in J_- : W(t) > c(u) + \lambda\}$ and $J_+(\lambda) := \{t \in J_+ : W(t) < c(u) - \lambda\}$.

Interestingly, positivity and integrability of the function e^W are sufficient to derive the desired estimate for I_2 , cf. also [32, Theorem 2.3]. In fact, with $\rho =$

 $\tau r^{2/\alpha}$ we have

$$e^{\lambda}\mu_{1}(J_{-}(\lambda)) = e^{\lambda}\mu_{1}(\{t \in J_{-} : e^{W(t)} > e^{c(u)}e^{\lambda}\}) = \int_{J_{-}(\lambda)} e^{\lambda} dt$$

$$\leq \int_{J_{-}(\lambda)} e^{W(t)-c(u)} dt \leq \int_{J_{-}} e^{W(t)-c(u)} dt$$

$$= \frac{g_{2-\alpha}(\eta\rho)}{(g_{1-\alpha} * e^{W})(\eta\rho)} \int_{0}^{\eta\rho} e^{W(t)} dt$$

$$\leq \frac{g_{2-\alpha}(\eta\rho)}{(g_{1-\alpha} * e^{W})(\eta\rho)} \cdot \frac{1}{g_{1-\alpha}(\eta\rho)} \int_{0}^{\eta\rho} g_{1-\alpha}(\eta\rho - t)e^{W(t)} dt$$

$$= \frac{\Gamma(1-\alpha)}{\Gamma(2-\alpha)} \eta\rho = \frac{\eta\tau r^{2/\alpha}}{1-\alpha},$$

and therefore

$$I_2 = \mu_1 \left(J_-(\lambda/2) \right) \mu_N(\delta B) \le \frac{2\eta \tau \delta^N}{(1-\alpha)\lambda} r^{2/\alpha} \mu_N(B), \quad \lambda > 0.$$
(3.44)

We come now to I_4 . For m > 0 define the function H_m on \mathbb{R} by $H_m(y) = y$, $y \le m$, and $H_m(y) = m + (y - m)/(y - m + 1)$, $y \ge m$. Then H_m is increasing, concave, and bounded above by m + 1. Further, we have $H_m \in C^1(\mathbb{R})$, and so by concavity

$$0 \le y H'_m(y) \le H_m(y) \le m+1, \quad y \ge 0.$$
(3.45)

Multiplying (3.40) by $e^W H'_m(e^W)$ and employing (3.45) as well as the fundamental identity (2.6), we infer that

$$\partial_t \left(g_{1-\alpha,n} * H_m(e^W) \right) + \frac{C_1}{r^2} H_m(e^W) \ge -S_n e^W H'_m(e^W), \quad \text{a.a. } t \in J.$$
 (3.46)

For $t \in J_+$ we shift the time by setting $s = t - \eta \tau r^{2/\alpha} = t - \eta \rho$ and put $\tilde{f}(s) = f(s + \eta \rho), s \in (0, (1 - \eta)\rho)$, for functions f defined on J_+ . By the time-shifting identity (3.2), (3.46) implies that for a.a. $s \in (0, (1 - \eta)\rho)$

$$\partial_{s}\left(g_{1-\alpha,n} \ast H_{m}\left(e^{\tilde{W}}\right)\right) + \frac{C_{1}}{r^{2}}H_{m}\left(e^{\tilde{W}}\right) \ge \Upsilon_{n,m}(s) - \tilde{S}_{n}e^{\tilde{W}}H_{m}'\left(e^{\tilde{W}}\right), \quad (3.47)$$

with the history term

$$\Upsilon_{n,m}(s) = \int_0^{\eta\rho} \left[-\dot{g}_{1-\alpha,n}(s+\eta\rho-\sigma) \right] H_m(e^{W(\sigma)}) \, d\sigma.$$

For $\theta \ge 0$ define the kernel $r_{\alpha,\theta} \in L_{1,\text{loc}}(\mathbb{R}_+)$ by means of

$$r_{\alpha,\,\theta}(t) + \theta(r_{\alpha,\,\theta} * g_{\alpha})(t) = g_{\alpha}(t), \quad t > 0.$$

Observe that $r_{\alpha,0} = g_{\alpha}$. Since g_{α} is completely monotone, $r_{\alpha,\theta}$ enjoys the same property (*cf.* [13, Chapter 5]), in particular $r_{\alpha,\theta}(s) > 0$ for all s > 0. Moreover, we have (see *e.g.* [32])

$$r_{\alpha,\theta}(s) = \Gamma(\alpha)g_{\alpha}(s) E_{\alpha,\alpha}(-\theta s^{\alpha}), \quad s > 0,$$

where $E_{\alpha,\beta}$ denotes the generalized Mittag-Leffler-function defined by

$$E_{\alpha,\beta}(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(n\alpha + \beta)}, \quad z \in \mathbb{C}.$$

We put $\theta = C_1/r^2$ and convolve (3.47) with $r_{\alpha,\theta}$. We have a.e. in $(0, (1 - \eta)\rho)$

$$\begin{aligned} r_{\alpha,\,\theta} * \partial_s \Big(g_{1-\alpha,n} * H_m(e^{\tilde{W}}) \Big) &= \partial_s \Big(r_{\alpha,\,\theta} * g_{1-\alpha,n} * H_m(e^{\tilde{W}}) \Big) \\ &= \partial_s \Big([g_\alpha - \theta(r_{\alpha,\,\theta} * g_\alpha)] * g_{1-\alpha,n} * H_m(e^{\tilde{W}}) \Big) \\ &= h_n * H_m(e^{\tilde{W}}) - \theta r_{\alpha,\,\theta} * h_n * H_m(e^{\tilde{W}}), \end{aligned}$$

and so we obtain a.e. in $(0, (1 - \eta)\rho)$

$$h_{n} * H_{m}(e^{\tilde{W}}) \geq r_{\alpha,\theta} * \Upsilon_{n,m} - r_{\alpha,\theta} * [\tilde{S}_{n}e^{\tilde{W}}H'_{m}(e^{\tilde{W}})] + \theta h_{n} * r_{\alpha,\theta} * H_{m}(e^{\tilde{W}}) - \theta r_{\alpha,\theta} * H_{m}(e^{\tilde{W}}).$$
(3.48)

Sending $n \to \infty$ and selecting an appropriate subsequence, if necessary, it follows that

$$H_m(e^W) \ge r_{\alpha,\,\theta} * \Upsilon_m, \qquad \text{a.a. } s \in (0,\,(1-\eta)\rho), \tag{3.49}$$

where

$$\Upsilon_m(s) = \int_0^{\eta\rho} \left[-\dot{g}_{1-\alpha}(s+\eta\rho-\sigma) \right] H_m(e^{W(\sigma)}) \, d\sigma.$$

Observe that for $s \in (0, (1 - \eta)\rho)$ we have

$$0 \le \theta s^{\alpha} \le \frac{C_1}{r^2} (1-\eta)^{\alpha} \left(\tau r^{2/\alpha}\right)^{\alpha} = C_1 (1-\eta)^{\alpha} \tau^{\alpha} =: \omega,$$

and thus by continuity and strict positivity of $E_{\alpha,\alpha}$ in $(-\infty, 0]$,

$$r_{\alpha,\theta}(s) \ge \Gamma(\alpha)g_{\alpha}(s) \min_{y \in [0,\omega]} E_{\alpha,\alpha}(-y) =: C_2(\alpha,\omega)\Gamma(\alpha)g_{\alpha}(s), \ s \in (0, (1-\eta)\rho).$$

We may then argue as in [32, Section 2.1] to obtain

$$H_m(e^{\tilde{W}(s)}) \ge C_2(\alpha, \omega) \frac{\alpha(s/[\eta\rho])^{\alpha}}{1 + (s/[\eta\rho])} (\eta\rho)^{\alpha-1} (g_{1-\alpha} * H_m(e^W))(\eta\rho),$$

a.a. $s \in (0, (1-\eta)\rho).$

Evidently, $H_m(y) \nearrow y$ as $m \to \infty$ for all $y \in \mathbb{R}$. Thus by sending $m \to \infty$ and applying Fatou's lemma we conclude that

$$e^{\tilde{W}(s)} \ge C_2(\alpha, \omega) \frac{\alpha(s/[\eta\rho])^{\alpha}}{1 + (s/[\eta\rho])} (\eta\rho)^{\alpha-1} (g_{1-\alpha} * e^W)(\eta\rho),$$

a.a. $s \in (0, (1-\eta)\rho).$ (3.50)

We then employ (3.50) to estimate as follows.

$$\begin{split} e^{\lambda}\mu_{1}\big(J_{+}(\lambda)\big) &= e^{\lambda}\mu_{1}\big(\{t \in J_{+} : e^{W(t)} < e^{c(u)}e^{-\lambda}\}\big) = \int_{J_{+}(\lambda)} e^{\lambda} dt \\ &\leq \int_{J_{+}(\lambda)} e^{c(u)-W(t)} dt \leq \int_{J_{+}} e^{c(u)-W(t)} dt \\ &= \frac{(g_{1-\alpha} * e^{W})(\eta\rho)}{g_{2-\alpha}(\eta\rho)} \int_{0}^{(1-\eta)\rho} e^{-\tilde{W}(s)} ds \\ &\leq \frac{C_{2}(\alpha, \omega)^{-1}(\eta\rho)^{1-\alpha}}{\alpha g_{2-\alpha}(\eta\rho)} \int_{0}^{(1-\eta)\rho} (1+s/\eta\rho)(s/\eta\rho)^{-\alpha} ds \\ &= \frac{\Gamma(2-\alpha)\eta\rho}{\alpha C_{2}(\alpha, \omega)} \int_{0}^{\frac{1-\eta}{\eta}} \sigma^{-\alpha}(1+\sigma) d\sigma = C_{3}(\alpha, \eta, \omega)\rho. \end{split}$$

Hence

$$I_4 = \mu_1 \left(J_+(\lambda/2) \right) \mu_N(\delta B) \le \frac{2C_3(\alpha, \eta, \omega)\delta^N}{\lambda} r^{2/\alpha} \mu_N(B), \quad \lambda > 0.$$
(3.51)

We come now to I_1 . Set $J_1(\lambda) = \{t \in J_- : c - W(t) + \lambda/2 \ge 0\}$ and $\Omega_t^-(\lambda) = \{x \in \delta B : w(t, x) > c + \lambda\}, t \in J_1(\lambda)$, where c = c(u) is given by (3.41). For $t \in J_1(\lambda)$, we have

$$w(t, x) - W(t) > c - W(t) + \lambda \ge \lambda/2, \quad x \in \Omega_t^-(\lambda),$$

and thus we deduce from (3.40) that a.e. in $J_1(\lambda)$

$$\frac{\nu}{16r^2\mu_N(B)} \mu_N\left(\Omega_t^-(\lambda)\right) \\
\leq \frac{1}{(c-W+\lambda)^2} \left(e^{-W}\partial_t(g_{1-\alpha,n} * e^W) + \frac{C_1}{r^2} + S_n\right).$$
(3.52)

Set $\chi(t, \lambda) = \mu_N(\Omega_t^-(\lambda))$, if $t \in J_1(\lambda)$, and $\chi(t, \lambda) = 0$ in case $t \in J_- \setminus J_1(\lambda)$. Let further $H(y) = (c - \log y + \lambda)^{-1}$, $0 < y \le y_* := e^{c+\lambda/2}$. Clearly, $H'(y) = (c - \log y + \lambda)^{-2}y^{-1}$ as well as

$$H''(y) = \frac{1}{(c - \log y + \lambda)^2 y^2} \left(\frac{2}{c - \log y + \lambda} - 1\right), \qquad 0 < y \le y_*,$$

which shows that H is concave in $(0, y_*]$ whenever $\lambda \ge 4$. We will assume this in what follows.

We next choose a C^1 extension \bar{H} of H on $(0, \infty)$ such that \bar{H} is concave, $0 \leq \bar{H}'(y) \leq \bar{H}'(y_*), y_* \leq y \leq 2y_*$, and $\bar{H}'(y) = 0, y \geq 2y_*$. Then

$$0 \le y\bar{H}'(y) \le \frac{2}{\lambda}, \quad y > 0.$$
(3.53)

In fact, for $y \in (0, y_*]$ we have

$$y\bar{H}'(y) = \frac{1}{(c-\log y+\lambda)^2} \le \frac{1}{(c-\log y_*+\lambda)^2} \le \frac{4}{\lambda^2} \le \frac{1}{\lambda},$$
 (3.54)

while in case $y \in [y_*, 2y_*]$ we may simply estimate

$$y\bar{H}'(y) \le 2y_*\bar{H}'(y_*) \le \frac{2}{\lambda}.$$

It is clear that \overline{H} is bounded above. There holds

$$\bar{H}(y) \le \frac{3}{\lambda}, \quad y > 0. \tag{3.55}$$

To see this, note that since \overline{H} is nondecreasing with $\overline{H}'(y) = 0$ for all $y \ge 2y_*$, the claim follows if the inequality is valid for all $y \in [y_*, 2y_*]$. For such y we have by (3.54) and by concavity of \overline{H}

$$\bar{H}(y) \le \bar{H}(y_*) + \bar{H}'(y_*)(y - y_*) \le \bar{H}(y_*) + y_*\bar{H}'(y_*) \le \frac{3}{\lambda}.$$

Observe also that

$$e^{W(t)}H'(e^{W(t)}) = \frac{1}{(c - W(t) + \lambda)^2}, \quad \text{a.a. } t \in J_1(\lambda).$$

Since $\bar{H}' \ge 0$, and $e^{-W} \partial_t (g_{1-\alpha,n} * e^W) + C_1 r^{-2} + S_n \ge 0$ on J_- by virtue of (3.40), we infer from (3.52) and (3.53) that

$$\frac{\nu}{16r^{2}\mu_{N}(B)}\chi(t,\lambda)$$

$$\leq e^{W}\bar{H}'(e^{W})\left(e^{-W}\partial_{t}(g_{1-\alpha,n}*e^{W})+\frac{C_{1}}{r^{2}}+S_{n}\right)$$

$$\leq \bar{H}'(e^{W})\partial_{t}(g_{1-\alpha,n}*e^{W})+\frac{2C_{1}}{\lambda r^{2}}+\frac{2|S_{n}(t)|}{\lambda}, \quad \text{a.a. } t \in J_{-}.$$
(3.56)

Since \overline{H} is concave, the fundamental identity (2.6) yields

$$\begin{split} \bar{H}'(e^W)\partial_t(g_{1-\alpha,n}*e^W) &\leq \partial_t \left(g_{1-\alpha,n}*\bar{H}(e^W)\right) + \left(-\bar{H}(e^W) + \bar{H}'(e^W)e^W\right)g_{1-\alpha,n} \\ &\leq \partial_t \left(g_{1-\alpha,n}*\bar{H}(e^W)\right) + \frac{2}{\lambda}g_{1-\alpha,n}, \quad \text{a.a. } t \in J_-, \end{split}$$

which, together with (3.56), gives a.e. in J_{-}

$$\frac{\nu}{16r^2\mu_N(B)}\chi(t,\lambda) \le \partial_t \Big((g_{1-\alpha,n} * \bar{H}(e^W) \Big) \\ + \frac{2}{\lambda} g_{1-\alpha,n} + \frac{2C_1}{\lambda r^2} + \frac{2|S_n(t)|}{\lambda}.$$
(3.57)

We then integrate (3.57) over $J_{-} = (0, \eta \rho)$ and employ (3.55) for the estimate

$$\left(g_{1-\alpha,n} * \bar{H}(e^W)\right)(\eta\rho) \leq \frac{3}{\lambda} \int_0^{\eta\rho} g_{1-\alpha,n}(t) dt$$

By sending $n \to \infty$, this leads to

$$\begin{split} &\int_{J_1(\lambda)} \mu_N \left(\Omega_t^-(\lambda) \right) dt = \int_0^{\eta \rho} \chi(t,\lambda) \, dt \le \frac{16r^2 \mu_N(B)}{\nu} \left(\frac{5}{\lambda} g_{2-\alpha}(\eta \rho) + \frac{2C_1 \eta \rho}{\lambda r^2} \right) \\ &= \frac{16r^{2/\alpha} \mu_N(B)}{\nu \lambda} \left(5g_{2-\alpha}(\eta \tau) + 2C_1 \eta \tau \right) =: C_4 \frac{r^{2/\alpha} \mu_N(B)}{\lambda}, \quad \lambda \ge 4. \end{split}$$

Hence with $C_5 = \max\{4\tau, C_4\}$ we find that

$$I_1 \le \frac{C_5 r^{2/\alpha} \mu_N(B)}{\lambda}, \quad \lambda > 0.$$
(3.58)

It remains to derive the desired estimate for I_3 . To this purpose we shift again the time by putting $s = t - \eta \rho$, and denote the corresponding transformed functions as above by $\tilde{W}, \tilde{w}, \ldots$ and so forth. Set further $\tilde{J}_+ := (0, (1 - \eta)\rho)$. By the time-shifting property (3.2) and by positivity of e^W , relation (3.40) then implies

$$\frac{\nu}{16r^2\mu_N(B)}\int_{\delta B}(\tilde{w}-\tilde{W})^2dx \le e^{-\tilde{W}}\partial_s(g_{1-\alpha,n}*e^{\tilde{W}}) + \frac{C_1}{r^2} + \tilde{S}_n(s),$$
(3.59)
a.a. $s \in \tilde{J}_+$.

Next, set $J_2(\lambda) = \{s \in \tilde{J}_+ : \tilde{W}(s) - c + \lambda/2 \ge 0\}$ and $\Omega_s^+(\lambda) = \{x \in \delta B : \tilde{W}(s, x) < c - \lambda\}, s \in J_2(\lambda)$. For $s \in J_2(\lambda)$, we have

$$\tilde{W}(s) - \tilde{w}(s, x) \ge \tilde{W}(s) - c + \lambda \ge \lambda/2, \quad x \in \Omega_s^+(\lambda),$$

and thus (3.59) yields that a.e. in $J_2(\lambda)$

$$\frac{\nu}{16r^2\mu_N(B)} \mu_N\left(\Omega_s^+(\lambda)\right) \\
\leq \frac{1}{(\tilde{W}-c+\lambda)^2} \left(e^{-\tilde{W}}\partial_s(g_{1-\alpha,n}*e^{\tilde{W}}) + \frac{C_1}{r^2} + \tilde{S}_n\right).$$
(3.60)

We proceed now similarly as above for the term I_1 . Set $\chi(s, \lambda) = \mu_N(\Omega_s^+(\lambda))$, if $s \in J_2(\lambda)$, and $\chi(s, \lambda) = 0$ in case $s \in \tilde{J}_+ \setminus J_1(\lambda)$. We consider this time the convex function $H(y) = (\log y - c + \lambda)^{-1}$ for $y \ge y_* := e^{c-\lambda/2}$ with derivative $H'(y) = -(\log y - c + \lambda)^{-2}y^{-1} < 0$. We define a C^1 extension \bar{H} of H on $[0, \infty)$ by means of

$$\bar{H}(y) = \begin{cases} H'(y_*)(y - y_*) + H(y_*) & 0 \le y < y_* \\ H(y) & y \ge y_*. \end{cases}$$

Evidently, $-\bar{H}$ is concave in $[0, \infty)$ and

$$0 \le -\bar{H}'(y)y \le \frac{1}{(\log y_* - c + \lambda)^2} \le \frac{1}{(\lambda/2)^2} \le \frac{4}{\lambda}, \quad y \ge 0, \ \lambda \ge 1.$$
(3.61)

We will assume $\lambda \ge 1$ in the subsequent lines.

Observe that

$$-e^{\tilde{W}(s)}H'(e^{\tilde{W}(s)}) = \frac{1}{(\tilde{W}(s) - c + \lambda)^2}, \quad \text{a.a. } s \in J_2(\lambda).$$

Since $-\bar{H}' \ge 0$, and $e^{-\tilde{W}} \partial_s(g_{1-\alpha,n} * e^{\tilde{W}}) + C_1 r^{-2} + \tilde{S}_n \ge 0$ on \tilde{J}_+ due to (3.59), it thus follows from (3.60) and (3.61) that

$$\frac{\nu}{16r^{2}\mu_{N}(B)} \chi(s,\lambda) \leq -e^{\tilde{W}}\bar{H}'(e^{\tilde{W}}) \left(e^{-\tilde{W}}\partial_{s}(g_{1-\alpha,n}*e^{\tilde{W}}) + \frac{C_{1}}{r^{2}} + \tilde{S}_{n}\right)$$

$$\leq -\bar{H}'(e^{\tilde{W}})\partial_{s}(g_{1-\alpha,n}*e^{\tilde{W}}) + \frac{4C_{1}}{\lambda r^{2}} + \frac{4|\tilde{S}_{n}(s)|}{\lambda}, \quad \text{a.a. } s \in \tilde{J}_{+}.$$
(3.62)

By concavity of $-\overline{H}$, the fundamental identity (2.6) provides the estimate

$$-\bar{H}'(e^{\tilde{W}})\partial_{s}(g_{1-\alpha,n}*e^{\tilde{W}}) \leq -\partial_{s}\left(g_{1-\alpha,n}*\bar{H}(e^{\tilde{W}})\right) + \left(\bar{H}(e^{\tilde{W}}) - \bar{H}'(e^{\tilde{W}})e^{\tilde{W}}\right)g_{1-\alpha,n}$$
$$\leq -\partial_{s}\left(g_{1-\alpha,n}*\bar{H}(e^{\tilde{W}})\right) + \bar{H}(0)g_{1-\alpha,n} \leq -\partial_{s}\left(g_{1-\alpha,n}*\bar{H}(e^{\tilde{W}})\right) + \frac{6}{\lambda}g_{1-\alpha,n},$$

a.e. in \tilde{J}_+ , which when combined with (3.62) leads to

$$\frac{\nu}{16r^2\mu_N(B)} \chi(s,\lambda) \le -\partial_s \Big(g_{1-\alpha,n} * \bar{H}(e^{\bar{W}}) \Big) + \frac{6}{\lambda} g_{1-\alpha,n} + \frac{4C_1}{\lambda r^2} + \frac{4|\tilde{S}_n(s)|}{\lambda}$$

for a.a. $s \in \tilde{J}_+$. We integrate this estimate over \tilde{J}_+ and send $n \to \infty$ to the result

$$\begin{split} \int_{J_2(\lambda)} \mu_N \big(\Omega_s^+(\lambda)\big) \, ds &= \int_0^{(1-\eta)\rho} \chi(s,\lambda) \, ds \\ &\leq \frac{16r^2 \mu_N(B)}{\nu} \left(\frac{6}{\lambda} g_{2-\alpha} \big((1-\eta)\rho\big) + \frac{4C_1(1-\eta)\rho}{\lambda r^2}\right) \\ &= \frac{16r^{2/\alpha} \mu_N(B)}{\nu \lambda} \big(6g_{2-\alpha} \big((1-\eta)\tau\big) + 4C_1(1-\eta)\tau\big) \\ &=: C_6 \frac{r^{2/\alpha} \mu_N(B)}{\lambda}, \quad \lambda \ge 1. \end{split}$$

Hence with $C_7 = \max{\tau, C_6}$ we obtain that

$$I_3 \le \frac{C_7 r^{2/\alpha} \mu_N(B)}{\lambda}, \quad \lambda > 0.$$
(3.63)

Finally, combining (3.42), (3.43), and (3.44), (3.51), (3.58), (3.63) establishes the theorem.

3.4. The final step

We are now in position to prove Theorem 1.1. Without loss of generality we may assume that $u \ge \varepsilon$ for some $\varepsilon > 0$; otherwise replace u by $u + \varepsilon$, which is a supersolution of (1.1) with $u_0 + \varepsilon$ instead of u_0 , and eventually let $\varepsilon \to 0+$.

For $0 < \sigma \leq 1$, we set $U_{\sigma} = (t_0 + (2 - \sigma)\tau r^{2/\alpha}, t_0 + 2\tau r^{2/\alpha}) \times \sigma B$ and $U'_{\sigma} = (t_0, t_0 + \sigma \tau r^{2/\alpha}) \times \sigma B$. Clearly, $Q_-(t_0, x_0, r) = U'_{\delta}$ and $Q_+(t_0, x_0, r) = U_{\delta}$. By Theorem 3.2,

$$\operatorname{ess\,sup}_{U_{\sigma'}} u^{-1} \leq \left(\frac{C\mu_{N+1}(U_1)^{-1}}{(\sigma - \sigma')^{\tau_0}}\right)^{1/\gamma} |u^{-1}|_{L_{\gamma}(U_{\sigma})}, \quad \delta \leq \sigma' < \sigma \leq 1, \ \gamma \in (0, 1].$$

Here $C = C(\nu, \Lambda, \delta, \tau, \alpha, N)$ and $\tau_0 = \tau_0(\alpha, N)$. This shows that the first hypothesis of Lemma 2.3 is satisfied by any positive constant multiple of u^{-1} with $\beta_0 = \infty$.

Consider now $f_1 = u^{-1}e^{c(u)}$ where c(u) is the constant from Theorem 3.4 with $K_- = U'_1$ and $K_+ = U_1$. Since $\log f_1 = c(u) - \log u$, we see from Theorem 3.4, estimate (3.32), that

$$\mu_{N+1}(\{(t, x) \in U_1 : \log f_1(t, x) > \lambda\}) \le M\mu_{N+1}(U_1)\lambda^{-1}, \quad \lambda > 0,$$

where $M = M(\nu, \Lambda, \delta, \tau, \eta, \alpha, N)$. Hence we may apply Lemma 2.3 with $\beta_0 = \infty$ to f_1 and the family U_{σ} ; thereby we obtain

$$\operatorname{ess\,sup}_{U_{\delta}} f_1 \le M_1$$

with $M_1 = M_1(\nu, \Lambda, \delta, \tau, \eta, \alpha, N)$. In terms of *u* this means that

$$e^{c(u)} \le M_1 \operatorname{essinf}_{U_{\delta}} u. \tag{3.64}$$

On the other hand, Theorem 3.3 yields

$$|u|_{L_{p}(U'_{\sigma'})} \leq \left(\frac{C\mu_{N+1}(U'_{1})^{-1}}{(\sigma - \sigma')^{\tau_{1}}}\right)^{1/\gamma - 1/p} |u|_{L_{\gamma}(U'_{\sigma})}, \quad \delta \leq \sigma' < \sigma \leq 1, \ 0 < \gamma \leq p/\tilde{\kappa}.$$

Here $C = C(v, \Lambda, \delta, \tau, \alpha, N, p)$ and $\tau_1 = \tau_1(\alpha, N)$. Thus the first hypothesis of Lemma 2.3 is satisfied by any positive constant multiple of u with $\beta_0 = p$ and $\eta = 1/\tilde{\kappa}$. Taking $f_2 = ue^{-c(u)}$ with c(u) from above, we have $\log f_2 = \log u - c(u)$ and so Theorem 3.4, estimate (3.31), gives

$$\mu_{N+1}(\{(t,x) \in U'_1 : \log f_2(t,x) > \lambda\}) \le M\mu_{N+1}(U'_1)\lambda^{-1}, \quad \lambda > 0,$$

where *M* is as above. Therefore we may again apply Lemma 2.3, this time to the function f_2 and the sets U'_{α} , and with $\beta_0 = p$ and $\eta = 1/\tilde{\kappa}$; we get

$$|f_2|_{L_p(U'_{\delta})} \le M_2 \mu_{N+1}(U'_1)^{1/p},$$

where $M_2 = M_2(\nu, \Lambda, \delta, \tau, \eta, \alpha, N, p)$. Rephrasing then yields

$$\mu_{N+1}(U_1')^{-1/p} |u|_{L_p(U_{\delta}')} \le M_2 e^{c(u)}.$$
(3.65)

Finally, we combine (3.64) and (3.65) to the result

$$\mu_{N+1}(U_1')^{-1/p} |u|_{L_p(U_{\delta}')} \le M_1 M_2 \, \operatorname*{essinf}_{U_{\delta}} u,$$

which proves the assertion.

4. Optimality of the exponent $\frac{2+N\alpha}{2+N\alpha-2\alpha}$ in the weak Harnack inequality

In this section we will show that the exponent $\frac{2+N\alpha}{2+N\alpha-2\alpha}$ in Theorem 1.1 is optimal. To this purpose consider the nonhomogeneous fractional diffusion equation

on \mathbb{R}^N

$$\partial_t^{\alpha} u - \Delta u = f, \quad t \in (0, T], \ x \in \mathbb{R}^N,$$
(4.1)

with initial condition

$$u(0, x) = 0, \quad x \in \mathbb{R}^n.$$
 (4.2)

Following [10], we say that a function $u \in C([0, T] \times \mathbb{R}^N) \cap C((0, T]; C^2(\mathbb{R}^N))$ with $g_{1-\alpha} * u \in C^1((0, T]; C(\mathbb{R}^N))$ is a classical solution of the problem (4.1), (4.2) if *u* satisfies (4.1) and (4.2). For any bounded continuous function *f* that is locally Hölder continuous in *x*, there exists a unique classical solution *u* of the problem (4.1), (4.2), and it is of the form

$$u(t,x) = \int_0^t \int_{\mathbb{R}^N} Y(t-\tau, x-y) f(\tau, y) \, dy \, d\tau,$$
(4.3)

where

$$Y(t,x) = c(N)|x|^{-N}t^{\alpha-1}H_{12}^{20}\left(\frac{1}{4}t^{-\alpha}|x|^2\Big|_{(N/2,1),(1,1)}^{(\alpha,\alpha)}\right),$$

cf. [10]. Here $H_{12}^{20}(z|_{(N/2,1),(1,1)}^{(\alpha,\alpha)})$ denotes a special H function (also termed Fox's H function), see [17, Section 1.12] and [10] for its definition. It is differentiable for z > 0, the asymptotic behaviour for $z \to \infty$ and $z \to +0$, respectively, is described in [10, formulae (3.9) and (3.14)]. It has been also proved in [10] that Y is nonnegative.

We choose a smooth and nonnegative approximation of unity $\{\phi_n(t, x)\}_{n \in \mathbb{N}}$ in $\mathbb{R}_+ \times \mathbb{R}^N$ such that each ϕ_n is bounded. Put $f = \phi_n$ in (4.1) and denote the corresponding classical solution of (4.1), (4.2) by u_n . Evidently, u_n is nonnegative and satisfies

$$\partial_t^{\alpha} u_n - \Delta u_n = \phi_n \ge 0, \quad t \in (0, T], \ x \in \mathbb{R}^N.$$

Hence u_n is a nonnegative supersolution of (4.1) with f = 0 for all $n \in \mathbb{N}$.

Suppose the weak Harnack inequality (1.3) holds for some $p \ge \frac{2+N\alpha}{2+N\alpha-2\alpha}$. Then, by taking $Q_- = (0, 1) \times B(0, 1)$ and $Q_+ = (2, 3) \times B(0, 1)$ it follows that

$$\left(\int_{Q_{-}} u_n^p \, d\mu_{N+1}\right)^{1/p} \le C \inf_{Q_{+}} u_n, \quad n \in \mathbb{N},\tag{4.4}$$

where the constant *C* is independent of *n*. Since $u_n \to Y$ in the distributional sense as $n \to \infty$, we have

$$\inf_{Q_{+}} u_{n} \leq \frac{1}{\mu_{N+1}(Q_{+})} \int_{Q_{+}} u_{n} d\mu_{N+1} \\
\leq 1 + \frac{1}{\mu_{N+1}(Q_{+})} \int_{Q_{+}} Y d\mu_{N+1} < \infty, \quad n \geq n_{0},$$

for a sufficiently large n_0 . On the other hand, the left-hand side of (4.4) cannot stay bounded, since $Y \notin L_p(Q_-)$ for $p \ge \frac{2+N\alpha}{2+N\alpha-2\alpha}$. In fact, writing $H_{12}^{20}(z) =$

 $H_{12}^{20}(z|_{(N/2,1),(1,1)}^{(\alpha,\alpha)})$ for short, we have

$$\begin{aligned} |Y|_{L_{p}(Q_{-})}^{p} &= \int_{0}^{1} \int_{B(0,1)} c(N)^{p} |x|^{-Np} t^{(\alpha-1)p} H_{12}^{20} (t^{-\alpha} |x|^{2}/4)^{p} \, dx \, dt \\ &= c_{1} \int_{0}^{1} \int_{0}^{1} r^{N-1-Np} t^{(\alpha-1)p} H_{12}^{20} (t^{-\alpha} r^{2}/4)^{p} \, dr \, dt \\ &= c_{1} \int_{0}^{1} \int_{0}^{t^{-\alpha/2}} (\rho t^{\alpha/2})^{N-1-Np} t^{(\alpha-1)p+\alpha/2} H_{12}^{20} (\rho^{2}/4)^{p} \, d\rho \, dt \\ &\geq c_{1} \int_{0}^{1} t^{\alpha(N-Np)/2+(\alpha-1)p} \, dt \int_{0}^{1} \rho^{N-1-Np} H_{12}^{20} (\rho^{2}/4)^{p} \, d\rho \\ &\geq c_{2} \int_{0}^{1} t^{\alpha(N-Np)/2+(\alpha-1)p} \, dt , \end{aligned}$$

with some positive constant c_2 . The last integral diverges for all $p \ge \frac{2+N\alpha}{2+N\alpha-2\alpha}$. Hence (4.4) yields a contradiction.

5. Applications of the weak Harnack inequality

The strong maximum principle for weak subsolutions of (1.1) may be easily derived as a consequence of the weak Harnack inequality.

Theorem 5.1. Let $\alpha \in (0, 1)$, T > 0, and $\Omega \subset \mathbb{R}^N$ be a bounded domain. Suppose the assumptions (H1)-(H3) are satisfied. Let $u \in Z_\alpha$ be a weak subsolution of (1.1) in Ω_T and assume that $0 \le \operatorname{ess\,sup}_{\Omega_T} u < \infty$ and that $\operatorname{ess\,sup}_{\Omega} u_0 \le \operatorname{ess\,sup}_{\Omega_T} u$. Then, if for some cylinder $Q = (t_0, t_0 + \tau r^{2/\alpha}) \times B(x_0, r) \subset \Omega_T$ with $t_0, \tau, r > 0$ and $\overline{B(x_0, r)} \subset \Omega$ we have

$$\operatorname{ess\,sup}_{Q} u = \operatorname{ess\,sup}_{\Omega_{T}} u, \tag{5.1}$$

the function *u* is constant on $(0, t_0) \times \Omega$.

Proof. Let $M = \operatorname{ess\,sup}_{\Omega_T} u$. Then v := M - u is a nonnegative weak supersolution of (1.1) with u_0 replaced by $v_0 := M - u_0 \ge 0$. For any $0 \le t_1 < t_1 + \eta r^{2/\alpha} < t_0$ the weak Harnack inequality with p = 1 applied to v yields an estimate of the form

$$r^{-(N+2/\alpha)} \int_{t_1}^{t_1+\eta r^{2/\alpha}} \int_{B(x_0,r)} (M-u) \, dx \, dt \le C \, \operatorname*{ess\,inf}_{\mathcal{Q}} (M-u) \, = \, 0.$$

This shows that u = M a.e. in $(0, t_0) \times B(x_0, r)$. As in the classical parabolic case (*cf.* [19]) the assertion now follows by a chaining argument.

We next apply the weak Harnack inequality to establish continuity at t = 0 for weak solutions.

Theorem 5.2. Let $\alpha \in (0, 1)$, T > 0, and $\Omega \subset \mathbb{R}^N$ be a bounded domain. Suppose the assumptions (H1) and (H2) are satisfied. Let $u \in Z_{\alpha}$ be a bounded weak solution of (1.1) in Ω_T with $u_0 = 0$. Then u is continuous at $(0, x_0)$ for all $x_0 \in \Omega$ and $\lim_{(t,x)\to(0,x_0)} u(t,x) = 0$. Moreover, letting $\eta > 0$ we have for any cylinder $Q(x_0, r_0) := (0, \eta r_0^{2/\alpha}) \times B(x_0, r_0) \subset \Omega_T$ and $r \in (0, r_0]$

$$\operatorname{ess osc}_{Q(x_0,r)} u \leq C \left(\frac{r}{r_0}\right)^{\delta} |u|_{L_{\infty}(\Omega_T)},$$
(5.2)

with

$$\operatorname{ess osc}_{Q(x_0,r)} = \operatorname{ess sup}_{Q(x_0,r)} - \operatorname{ess inf}_{Q(x_0,r)}$$

and constants $C = Cf(v, \Lambda, \eta, \alpha, N) > 0$ and $\delta = \delta(v, \Lambda, \eta, \alpha, N) \in (0, 1)$.

Proof. Let $u \in Z_{\alpha}$ be a bounded weak solution of (1.1) in Ω_T with $u_0 = 0$. Set u(t, x) = 0 and A(t, x) = Id for t < 0 and $x \in \Omega$. For $T_0 > 0$ we shift the time by setting $s = t + T_0$ and put $\tilde{f}(s) = f(s - T_0), s \in (0, T + T_0)$, for functions f defined on $(-T_0, T)$. Since $Du(t, \cdot) = 0$ for t < 0 and

$$\partial_t (g_{1-\alpha,n} * u)(t,x) = \partial_t \int_{-T_0}^t g_{1-\alpha,n}(t-\tau) u(\tau,x) d\tau = \partial_s (g_{1-\alpha,n} * \tilde{u})(s,x),$$

the function \tilde{u} is a bounded weak solution of

$$\partial_s^{\alpha} \tilde{u} - \operatorname{div} \left(\tilde{A}(s, x) D \tilde{u} \right) = 0, \quad s \in (0, T + T_0), \, x \in \Omega.$$

Next, assuming $r \in (0, r_0/2]$ we introduce the cylinders

$$Q_*(x_0, r) = \left(-\eta r^{2/\alpha}, \eta r^{2/\alpha}\right) \times B(x_0, r), Q_-(x_0, r) = \left(-\eta (2r)^{2/\alpha}, -\eta (3r/2)^{2/\alpha}\right) \times B(x_0, r),$$

and denote by $\tilde{Q}_*(x_0, r)$ resp. $\tilde{Q}_-(x_0, r)$ the corresponding cylinders in the (s, x) coordinate system. Let us write $M_i = \operatorname{ess} \sup_{\tilde{Q}_*(x_0,ir)} \tilde{u}$ and $m_i = \operatorname{ess} \inf_{\tilde{Q}_*(x_0,ir)} \tilde{u}$ for i = 1, 2. Choosing $T_0 \ge \eta (2r)^{2/\alpha}$, we may apply Theorem 1.1 with p = 1 to the functions $M_2 - \tilde{u}, \tilde{u} - m_2$, which are nonnegative in $(0, \eta (2r)^{2/\alpha} + T_0) \times B(x_0, 2r)$, thereby obtaining

$$r^{-N+2/\alpha} \int_{\tilde{Q}_{-}(x_{0},r)} (M_{2}-\tilde{u}) d\mu_{N+1} \leq C(M_{2}-M_{1}),$$

$$r^{-N+2/\alpha} \int_{\tilde{Q}_{-}(x_{0},r)} (\tilde{u}-m_{2}) d\mu_{N+1} \leq C(m_{1}-m_{2}),$$

where C > 1 is a constant independent of u and r. By addition, it follows that

$$M_2 - m_2 \le C(M_2 - m_2 + m_1 - M_1).$$

Writing $\omega(x_0, r) = \operatorname{ess\,sup}_{\tilde{Q}_*(x_0, ir)} \tilde{u} - \operatorname{ess\,inf}_{\tilde{Q}_*(x_0, ir)} \tilde{u}$, this yields

$$\omega(x_0, r) \le \theta \omega(x_0, 2r), \quad r \le r_0/2, \tag{5.3}$$

where $\theta = 1 - C^{-1} \in (0, 1)$. Iterating (5.3) as in the proof of [11, Lemma 8.23] we obtain

$$\omega(x_0, r) \le \frac{1}{\theta} \left(\frac{r}{r_0} \right)^{\log \theta / \log(1/2)} \omega(x_0, r_0), \quad r \le r_0.$$

The estimate (5.2) then follows by transforming back to the function u and using that u = 0 for negative times. In particular, we also see that u is continuous at $(0, x_0)$ for all $x_0 \in \Omega$ and that $\lim_{(t,x)\to(0,x_0)} u(t, x) = 0$.

The last application is a uniqueness theorem for global bounded weak solutions. We say that a function u on $\mathbb{R}_+ \times \mathbb{R}^N$ is a *global weak solution* of

$$\partial_t^{\alpha} u - \operatorname{div} \left(A(t, x) D u \right) = 0, \tag{5.4}$$

if it is a weak solution of (5.4) in $(0, T) \times B(0, r)$ for all T > 0 and r > 0.

Corollary 5.3. Let $\alpha \in (0, 1)$. Assume that $A \in L_{\infty}(\mathbb{R}_+ \times \mathbb{R}^N; \mathbb{R}^{N \times N})$ and that there exists $\nu > 0$ such that

$$(A(t, x)\xi|\xi) \ge \nu|\xi|^2$$
, for a.a. $(t, x) \in \mathbb{R}_+ \times \mathbb{R}^N$, and all $\xi \in \mathbb{R}^N$.

Suppose that u is a global bounded weak solution of (5.4). Then u = 0 a.e. on $\mathbb{R}_+ \times \mathbb{R}^N$.

Proof. For r > 0 and $x_0 = 0$ it follows from the proof of Theorem 5.2 that

$$\omega(0,r) \le \theta \omega(0,2r), \quad r > 0, \tag{5.5}$$

where $\theta \in (0, 1)$ is independent of r and u. By induction, (5.5) yields

$$\omega(0,r) \leq \theta^n \omega(0,2^n r) \leq 2\theta^n |u|_{L_{\infty}(\mathbb{R}_+ \times \mathbb{R}^N)}, \quad r > 0, n \in \mathbb{N}.$$

Sending $n \to \infty$ shows that *u* is constant. The claim then follows by Theorem 5.2.

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